

# Sensitivity analysis for the Euler equations in Lagrangian coordinates

**LMV**  
Laboratoire de mathématiques  
de Versailles - CNRS UMR 8100



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# Outline of the talk

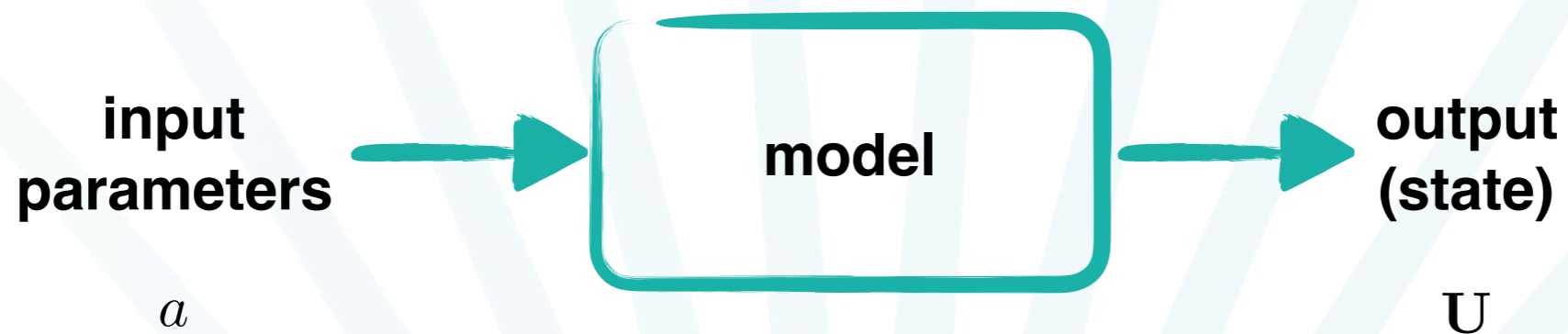
- ▶ Sensitivity analysis
- ▶ Sensitivity analysis for hyperbolic equations
- ▶ Euler equations in barotropic conditions ( $p$ -system)
- ▶ Classical numerical schemes and results
- ▶ Anti-diffusive numerical scheme and results



# **Sensitivity Analysis**

# Sensitivity Analysis

Sensitivity analysis: study of how changes in the **inputs** of a model affect the **outputs**



**Sensitivity:**  $\frac{\partial U}{\partial a} = U_a$

# Applications

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- ▶ Optimization

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► Optimization

Problem:  $\min_{a \in \mathcal{A}} J(\mathbf{U})$ , where  $J(\mathbf{U}) = \frac{1}{2}b(\mathbf{U}, \mathbf{U})$  and  $b$  is bilinear.

Classical optimization techniques call for the differentiation of the cost function:

$$\frac{\partial J(\mathbf{U})}{\partial a} = b(\mathbf{U}, \mathbf{U}_a)$$

# Applications

- ▶ Optimization
- ▶ Quick evaluation of close solutions



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$$\mathbf{U}(a + \delta a) = \mathbf{U}(a) + \delta a \mathbf{U}_a(a) + o(\delta a^2)$$



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- ▶ Quick evaluation of close solutions
- ▶ Uncertainty quantification

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First order estimates

$$\begin{array}{l} \mu \\ \sigma^2 \end{array} \quad \begin{array}{l} \mathbf{U}(\mu_a) \\ \mathbf{U}_a(\mu_a)^T \mathbf{U}_a(\mu_a) \sigma_a^2 \end{array}$$

# Standard techniques of sensitivity analysis

Standard techniques of sensitivity analysis call for the differentiation of the state system:

$$\begin{cases} \partial_t \mathbf{U} + \nabla \mathbf{F}(\mathbf{U}) = 0 & \Omega \times (0, T), \\ \mathbf{U}(x, 0) = \mathbf{g}(x) & \Omega, \end{cases}$$

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This can be done under **hypotheses of regularity** of the state  $\mathbf{U}$ .



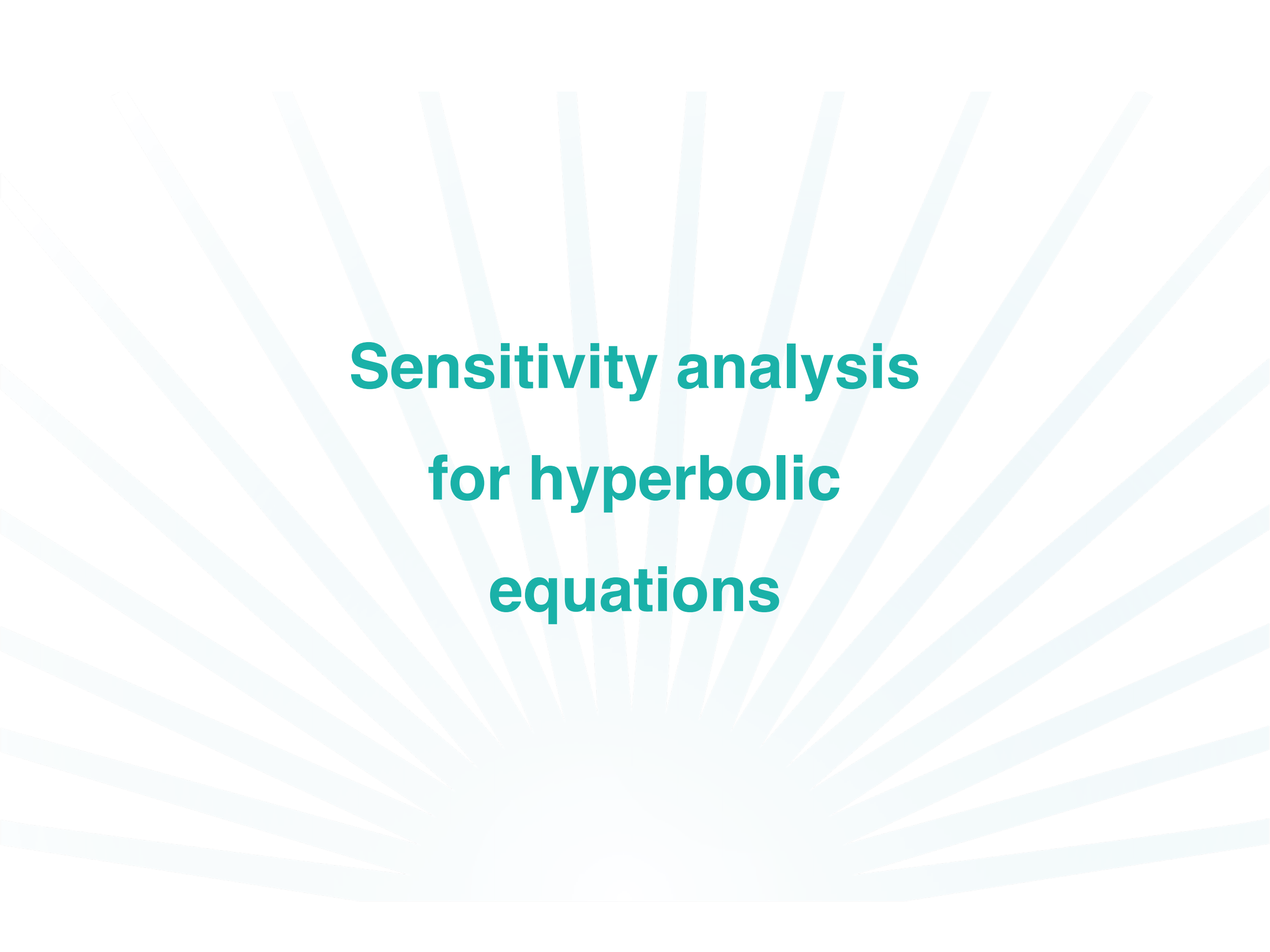
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If these techniques are applied to hyperbolic equations, **Dirac delta functions** will appear in the sensitivity.



**Sensitivity analysis  
for hyperbolic  
equations**

# Definition of the source term

In order to have a sensitivity system which is valid also when the state is discontinuous, we add a correction term:

$$\begin{cases} \partial_t \mathbf{U}_a + \partial_x \mathbf{F}_a(\mathbf{U}, \mathbf{U}_a) = \mathbf{S} & \Omega \times (0, T), \\ \mathbf{U}_a(x, 0) = \mathbf{g}_a(x) & \Omega, \end{cases}$$

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$$\mathbf{S} = \sum_{k=1}^{N_s} \rho_k \delta(x - x_{k,s}(t)),$$

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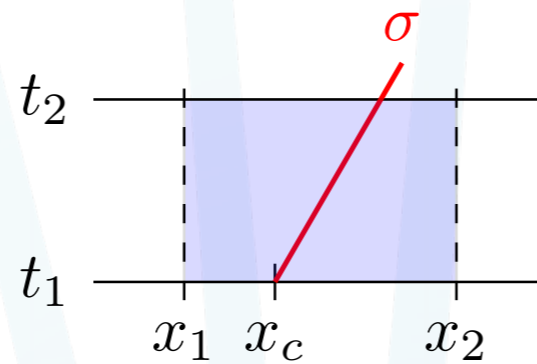
position of the k-th discontinuity

amplitude of the k-th correction  
(to be computed)



# Definition of the source term

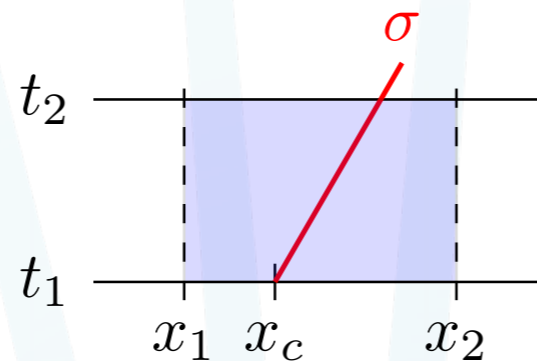
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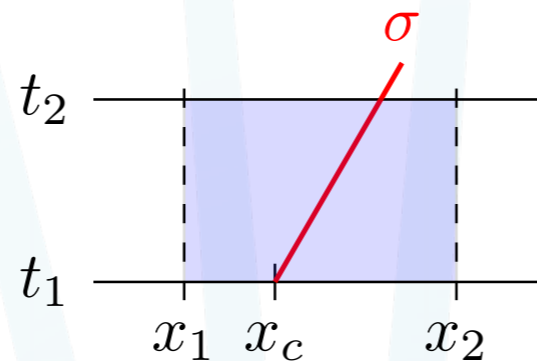


By integrating the sensitivity equations with the source term on the control volume, one has:

$$\rho = (\mathbf{U}_a^- - \mathbf{U}_a^+) \sigma + \mathbf{F}_a^+ - \mathbf{F}_a^-$$

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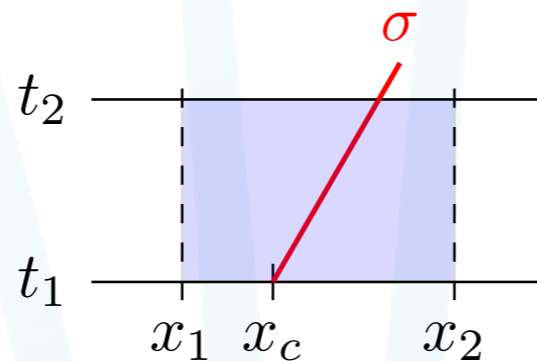
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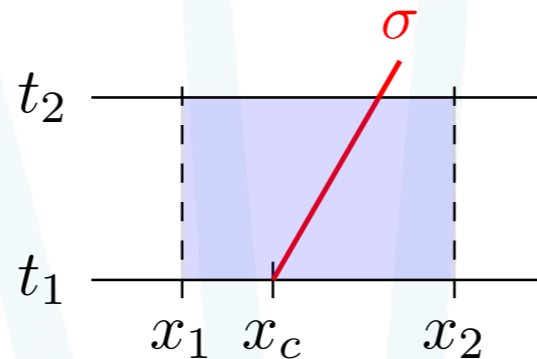
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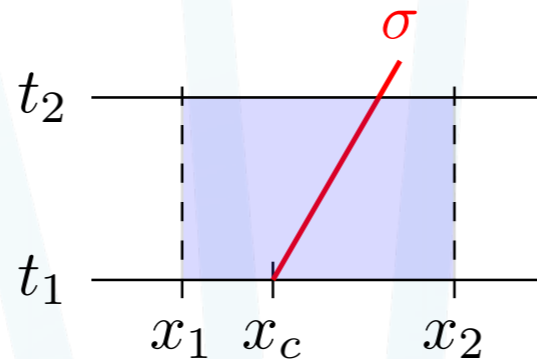
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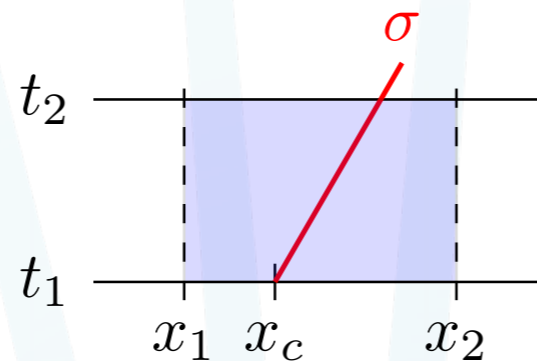
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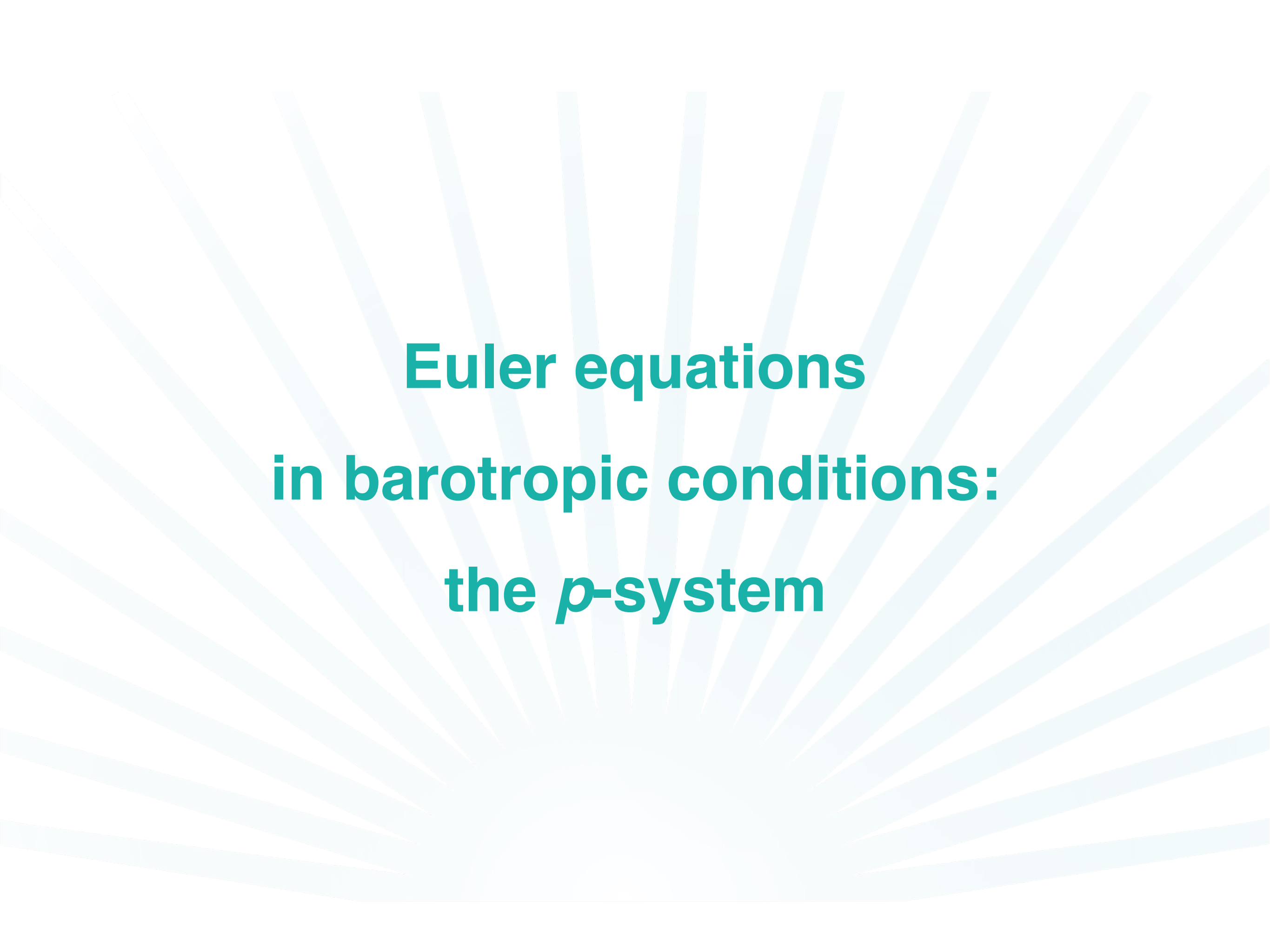
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**A shock detector is necessary**



**Euler equations  
in barotropic conditions:  
the  $p$ -system**



# The Riemann problem for the $p$ -system

The  $p$ -system writes:

$$\begin{cases} \partial_t \tau - \partial_x u = 0, \\ \partial_t u + \partial_x p(\tau) = 0. \end{cases}$$

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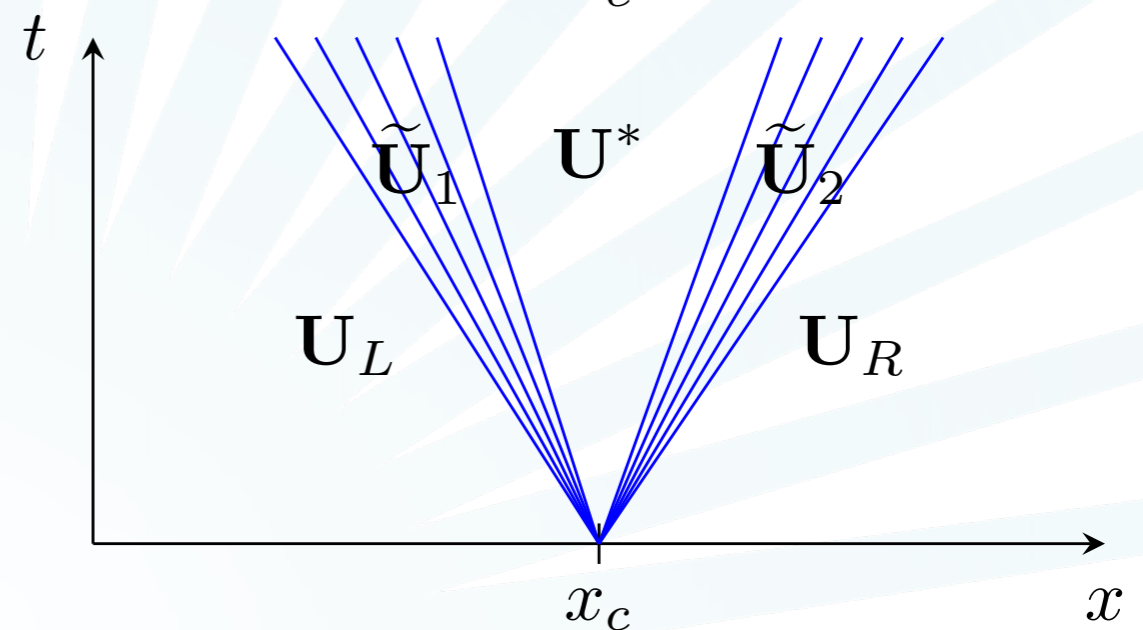
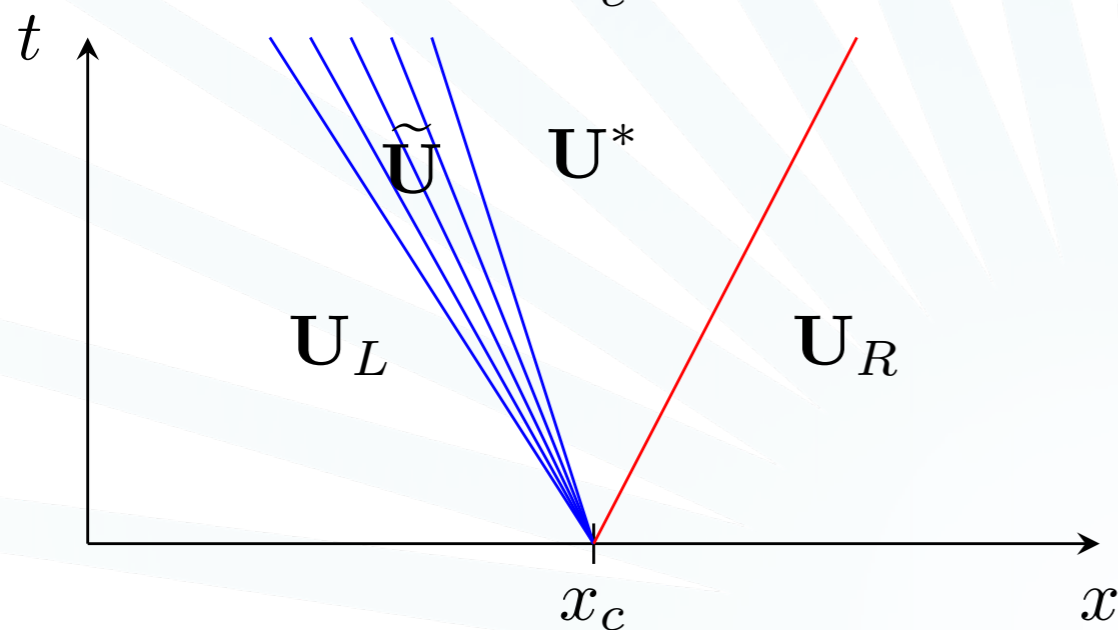
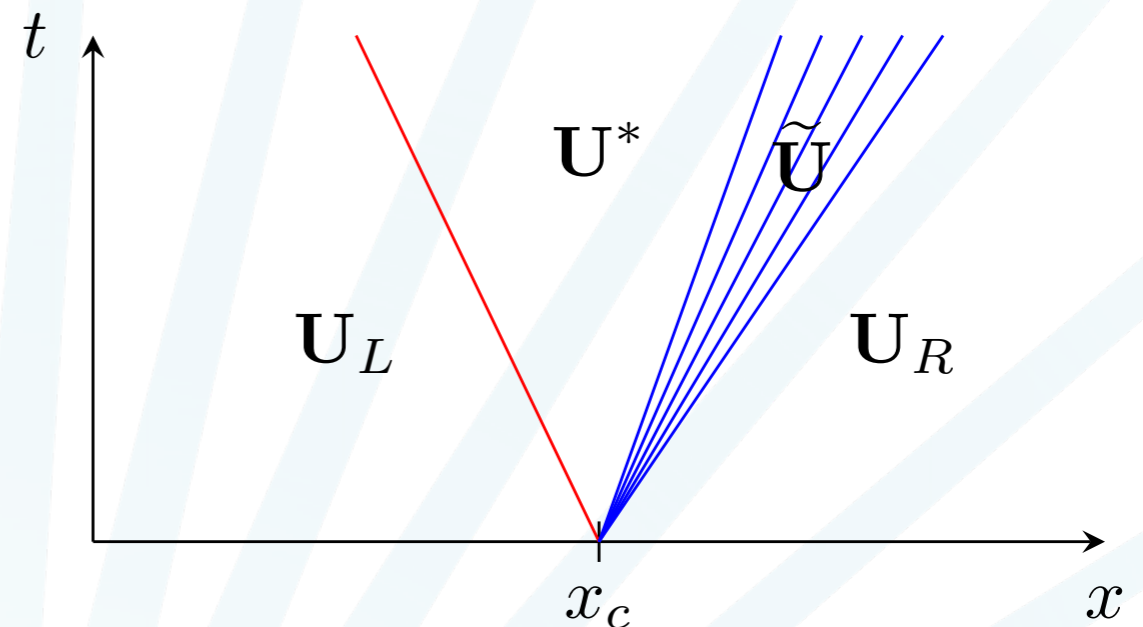
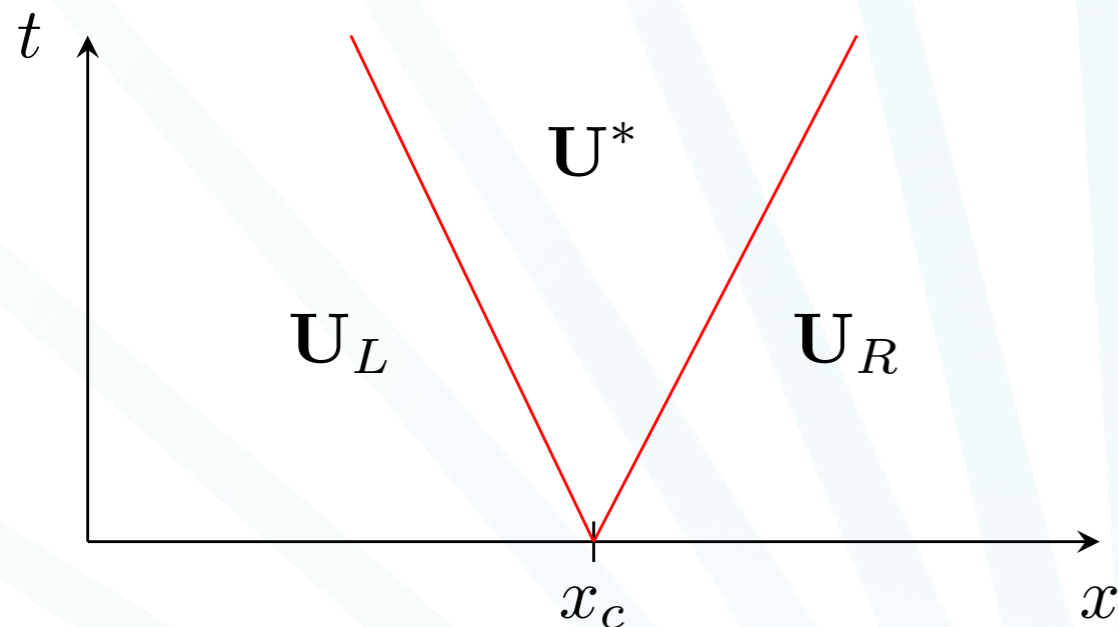
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$$g_1(\tau; \mathbf{U}_L) = \begin{cases} u_L - \sqrt{-(\tau^{-\gamma} - \tau_L^{-\gamma})(\tau - \tau_L)} & \text{if } \tau \leq \tau_L, \\ u_L + \frac{2\sqrt{\gamma}}{1-\gamma} (\tau^{\frac{1-\gamma}{2}} - \tau_L^{\frac{1-\gamma}{2}}) & \text{if } \tau > \tau_L. \end{cases}$$

$$g_2(\tau; \mathbf{U}_R) = \begin{cases} u_R + \sqrt{-(\tau^{-\gamma} - \tau_R^{-\gamma})(\tau - \tau_R)} & \text{if } \tau \leq \tau_R, \\ u_R + \frac{2\sqrt{\gamma}}{1-\gamma} (\tau_R^{\frac{1-\gamma}{2}} - \tau^{\frac{1-\gamma}{2}}) & \text{if } \tau > \tau_R. \end{cases}$$

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$$\tilde{\tau}_1 = \left( \frac{\gamma t^2}{(x - x_c)^2} \right)^{\frac{1}{\gamma+1}} \quad \tilde{u}_1 = u_L + \frac{2\sqrt{\gamma}}{1-\gamma} \left( \tilde{\tau}_1^{\frac{1-\gamma}{2}} - \tau_L^{\frac{1-\gamma}{2}} \right)$$

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# The sensitivity system

The sensitivity system writes:

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$$\begin{aligned} u_a^* &= g_1'(\tau^*; \mathbf{U}_L) \tau_a^* + \frac{\partial g_1}{\partial \tau_L}(\tau^*; \mathbf{U}_L) \tau_{a,L} + \frac{\partial g_1}{\partial u_L}(\tau^*; \mathbf{U}_L) u_{a,L} = \\ &= g_2'(\tau^*; \mathbf{U}_R) \tau_a^* + \frac{\partial g_2}{\partial \tau_R}(\tau^*; \mathbf{U}_R) \tau_{a,R} + \frac{\partial g_2}{\partial u_R}(\tau^*; \mathbf{U}_R) u_{a,R}, \end{aligned}$$

$$\tau_a^* = \frac{\frac{\partial g_2}{\partial \tau_R}(\tau^*; \mathbf{U}_R) \tau_{a,R} + \frac{\partial g_2}{\partial u_R}(\tau^*; \mathbf{U}_R) u_{a,R} - \frac{\partial g_1}{\partial \tau_L}(\tau^*; \mathbf{U}_L) \tau_{a,L} - \frac{\partial g_1}{\partial u_L}(\tau^*; \mathbf{U}_L) u_{a,L}}{g_1'(\tau^*; \mathbf{U}_L) - g_2'(\tau^*; \mathbf{U}_R)}$$

$$\tilde{\tau}_{a,1} = 0$$

$$\tilde{u}_{a,1} = \frac{\partial \tilde{u}_1}{\partial a} = u_{a,L} - \sqrt{\gamma} \tau_L^{-\frac{1+\gamma}{2}} \tau_{a,L}$$

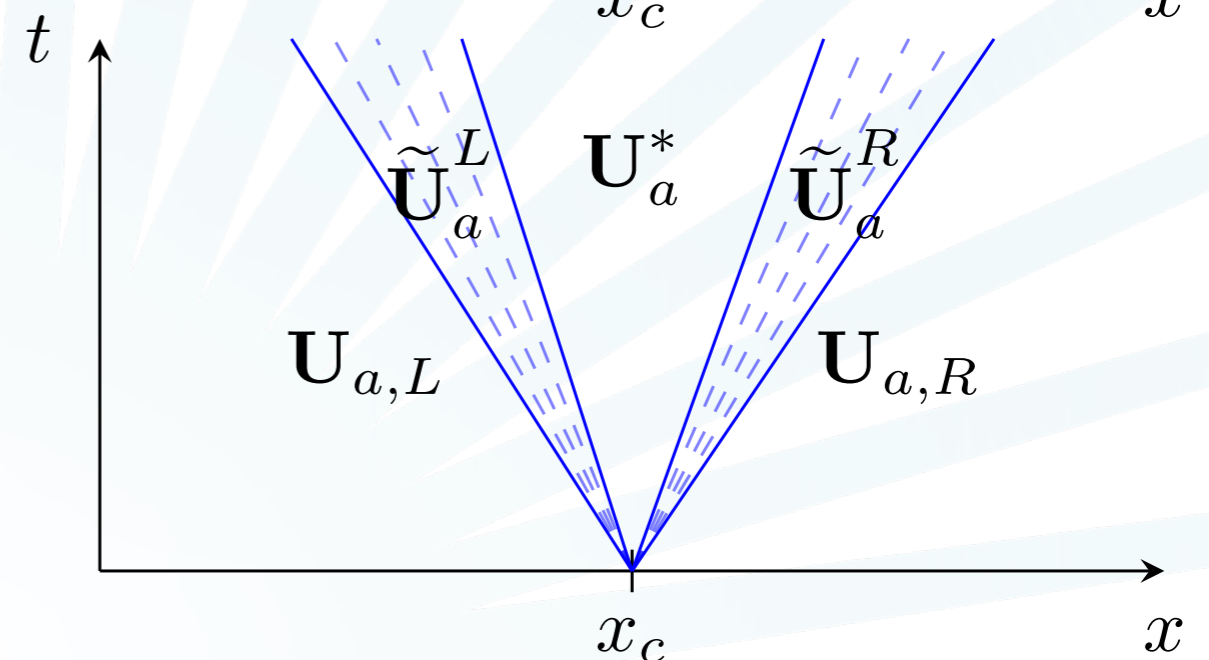
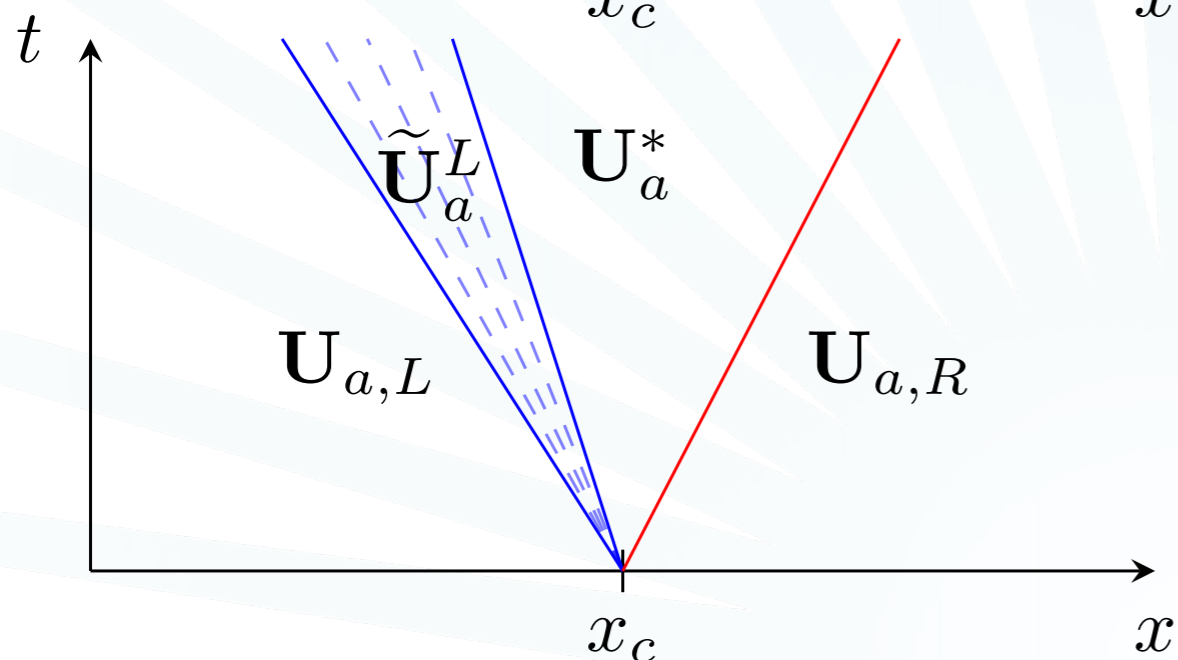
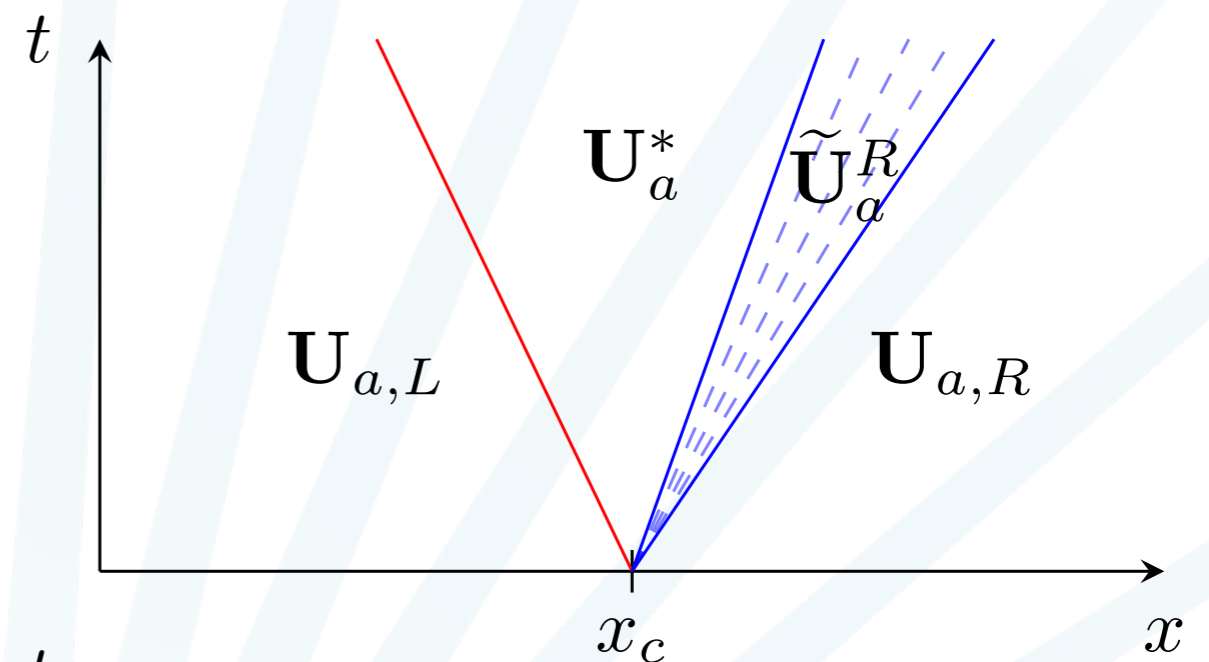
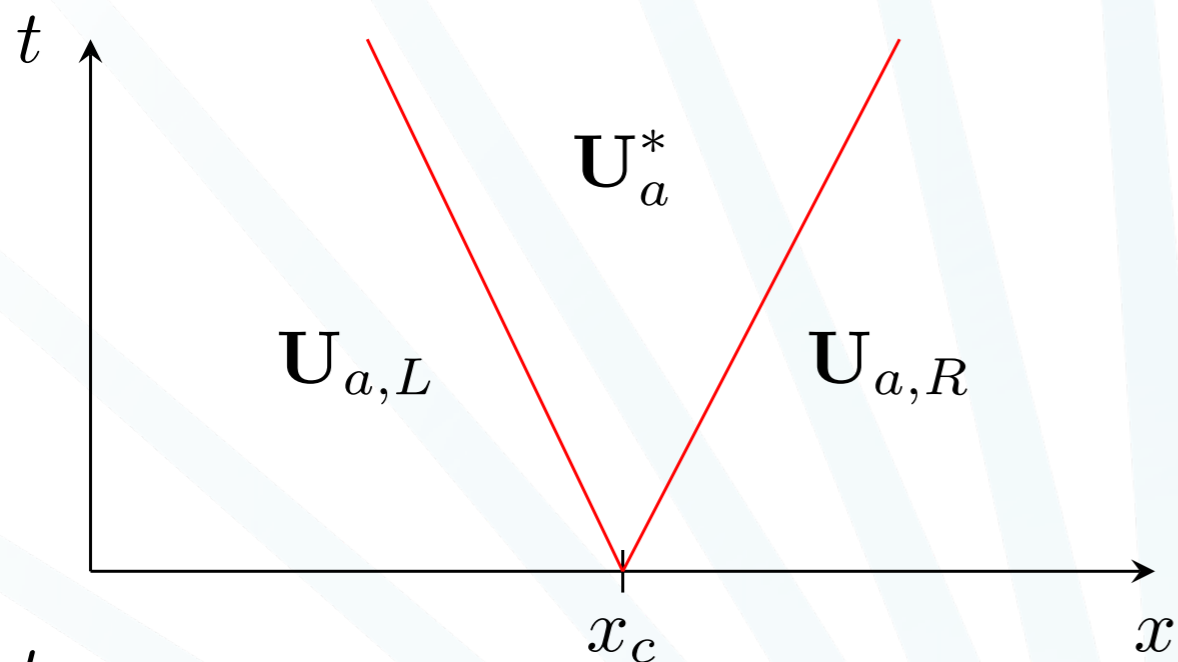
$$\tilde{\tau}_{a,2} = 0$$

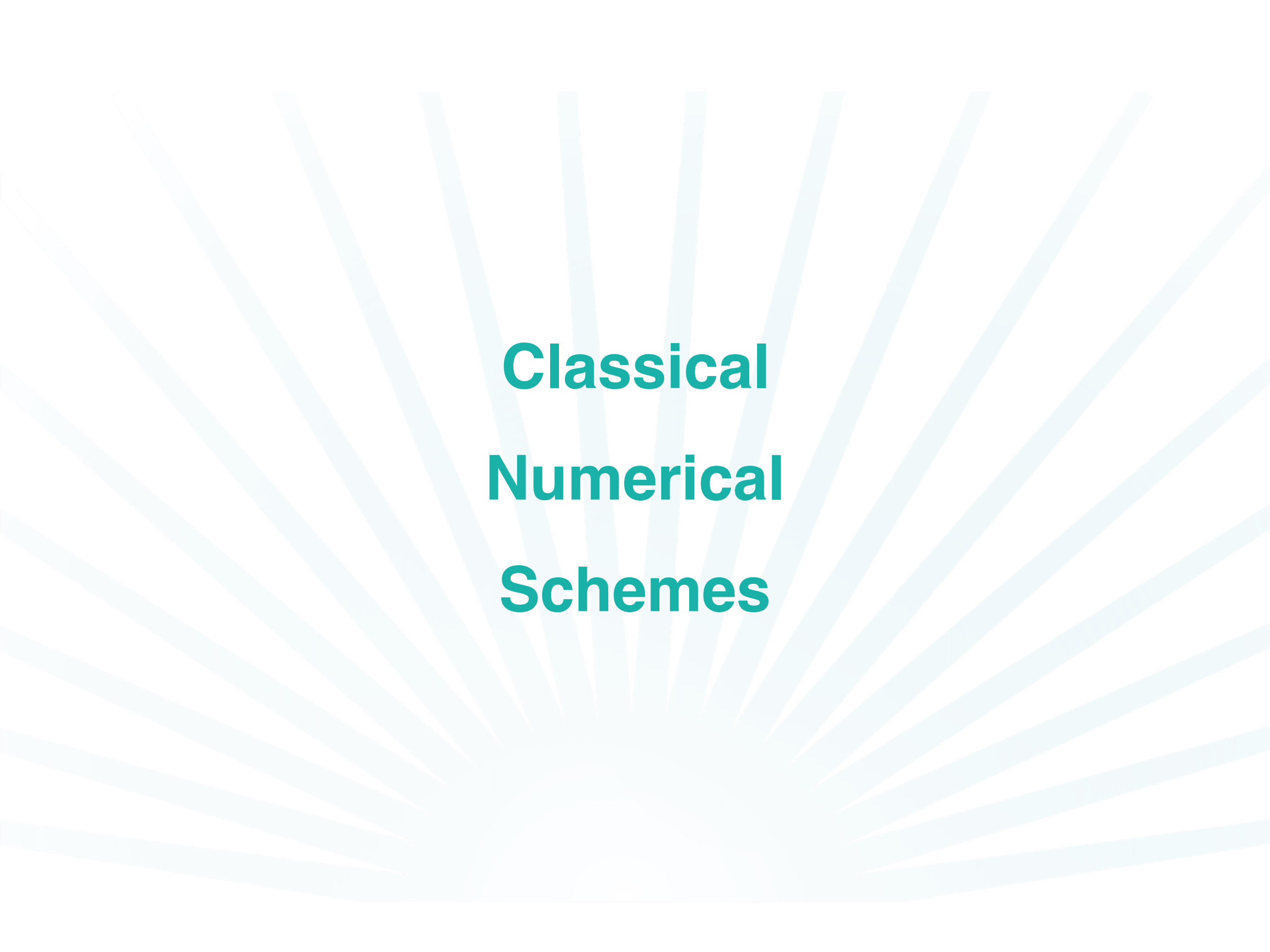
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**Classical  
Numerical  
Schemes**

# Classical numerical schemes





# Classical numerical schemes

- ▶ Exact Godunov-type scheme

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► Exact Godunov-type scheme

State: 
$$\mathbf{U}_j^{n+1} = \mathbf{U}_j^n - \frac{\Delta t}{\Delta x} (\mathbf{F}(\mathbf{U}_{j+1/2}^*) - \mathbf{F}(\mathbf{U}_{j-1/2}^*))$$

Sensitivity: non conservative, but composed **only** by **discontinuities**  
**direct average**

# Classical numerical schemes

- ▶ Exact Godunov-type scheme
- ▶ First order Roe-type scheme

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- ▶ Exact Godunov-type scheme
- ▶ First order Roe-type scheme

State: 
$$\mathbf{U}_j^{n+1} = \mathbf{U}_j^n + \frac{\Delta t}{\Delta x} (\lambda_{j-1/2}^{ROE} (\mathbf{U}_{j-1/2}^* - \mathbf{U}_j^n) + \lambda_{j+1/2}^{ROE} (\mathbf{U}_{j+1/2}^* - \mathbf{U}_j^n))$$

Sensitivity: the source term is encompassed in the definition of  $\mathbf{U}_{a,j-1/2}^*$  :

$$\begin{aligned} \mathbf{U}_{a,j-1/2}^* &= \frac{1}{2} (\mathbf{U}_{a,j-1}^n + \mathbf{U}_{a,j}^n) - \frac{\mathbf{F}_a(\mathbf{U}_j^n, \mathbf{U}_{a,j}^n) - \mathbf{F}_a(\mathbf{U}_{j-1}^n, \mathbf{U}_{a,j}^n)}{2\lambda_{j-1/2}^{ROE}} \\ &+ \frac{\lambda_{a,j-1/2}^{ROE}}{2\lambda_{j-1/2}^{ROE}} \left( (\mathbf{U}_{j-1}^n - \mathbf{U}_{j-1/2}^*) d_{1,j-1} + (\mathbf{U}_j^n - \mathbf{U}_{j-1/2}^*) d_{2,j} \right). \end{aligned}$$

# Classical numerical schemes

- ▶ Exact Godunov-type scheme
- ▶ First order Roe-type scheme

State: 
$$\mathbf{U}_j^{n+1} = \mathbf{U}_j^n + \frac{\Delta t}{\Delta x} (\lambda_{j-1/2}^{ROE} (\mathbf{U}_{j-1/2}^* - \mathbf{U}_j^n) + \lambda_{j+1/2}^{ROE} (\mathbf{U}_{j+1/2}^* - \mathbf{U}_j^n))$$

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Shock detectors

# Classical numerical schemes

- ▶ Exact Godunov-type scheme
- ▶ First order Roe-type scheme
- ▶ Second order Roe-type scheme



# Classical numerical schemes

- ▶ Exact Godunov-type scheme
- ▶ First order Roe-type scheme
- ▶ Second order Roe-type scheme

Time discretisation: two-step Runge-Kutta method

# Classical numerical schemes

- ▶ Exact Godunov-type scheme
- ▶ First order Roe-type scheme
- ▶ Second order Roe-type scheme

Time discretisation: two-step Runge-Kutta method

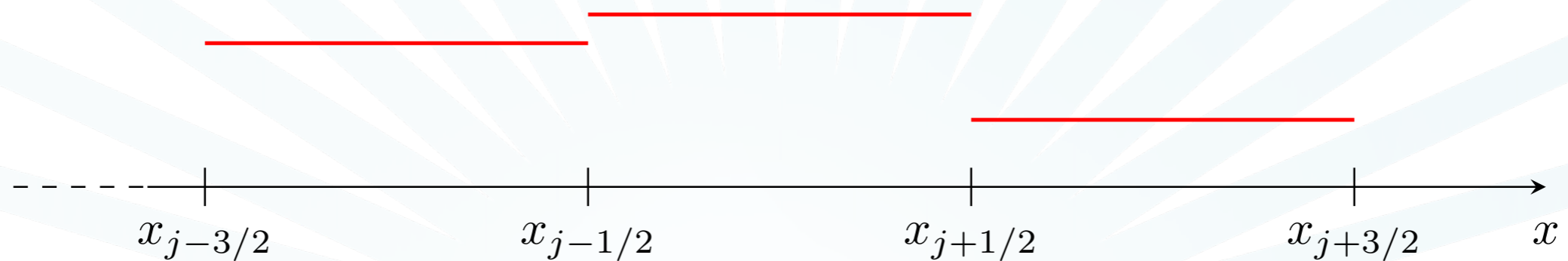
Space discretisation: MUSCL-type scheme

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Time discretisation: two-step Runge-Kutta method

Space discretisation: MUSCL-type scheme

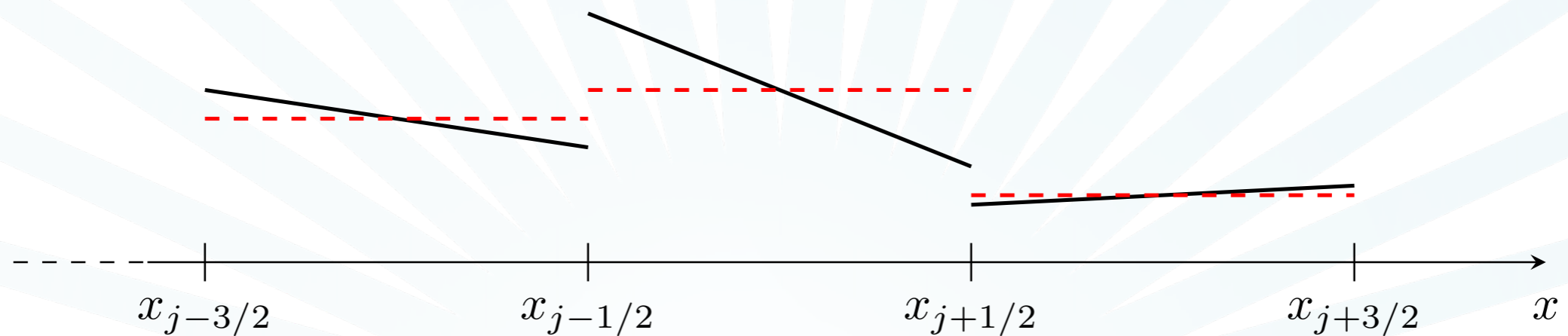


# Classical numerical schemes

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Time discretisation: two-step Runge-Kutta method

Space discretisation: MUSCL-type scheme

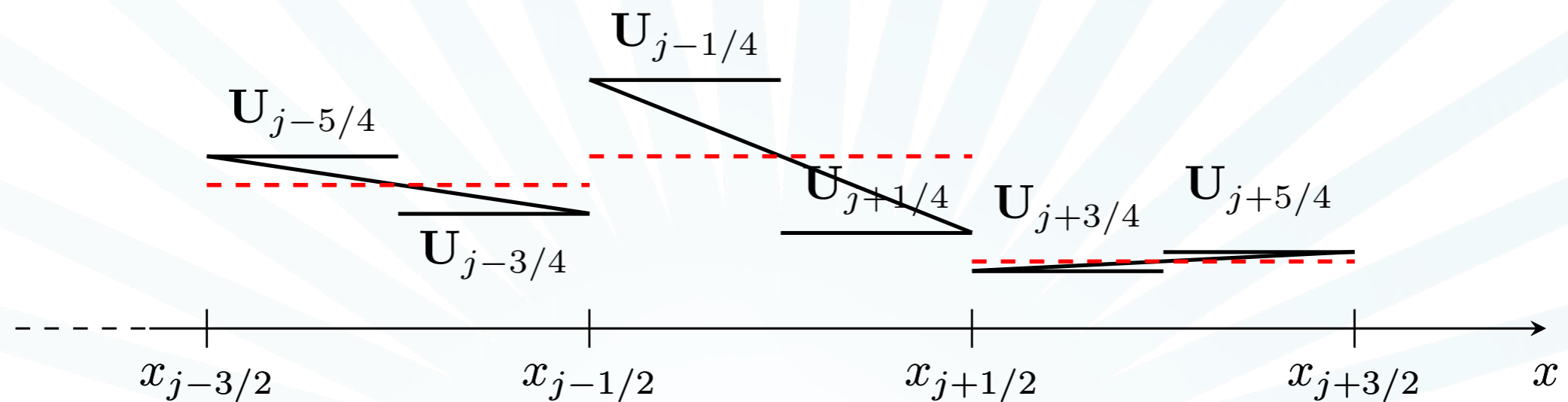


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- ▶ Exact Godunov-type scheme
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Time discretisation: two-step Runge-Kutta method

Space discretisation: MUSCL-type scheme

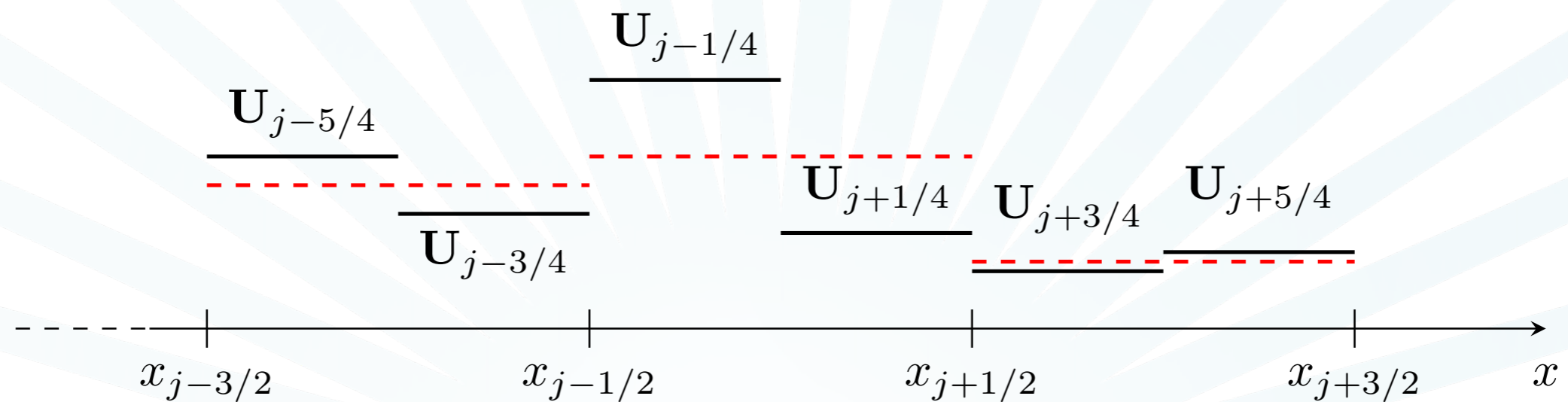


# Classical numerical schemes

- ▶ Exact Godunov-type scheme
- ▶ First order Roe-type scheme
- ▶ Second order Roe-type scheme

Time discretisation: two-step Runge-Kutta method

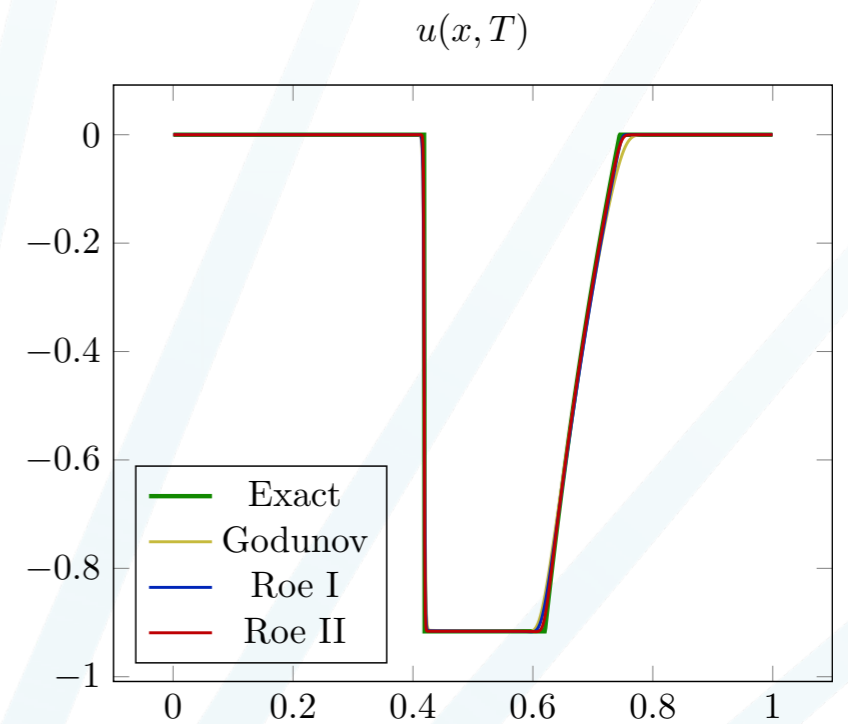
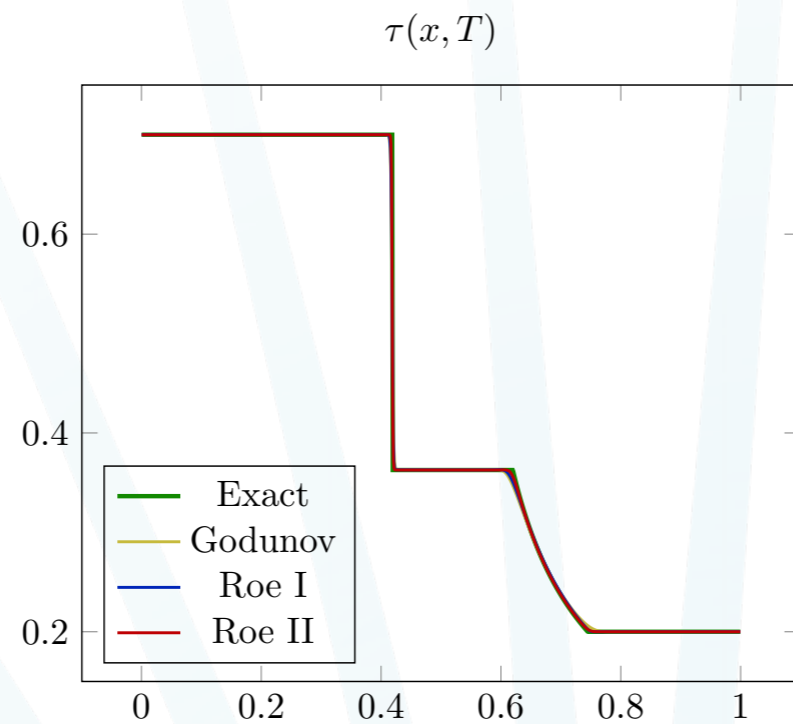
Space discretisation: MUSCL-type scheme





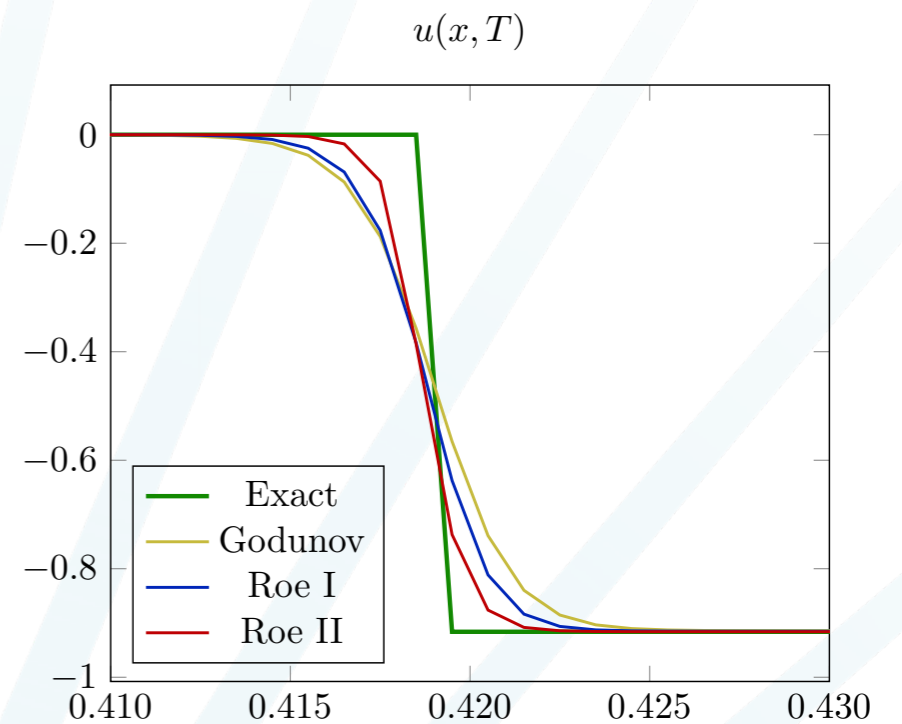
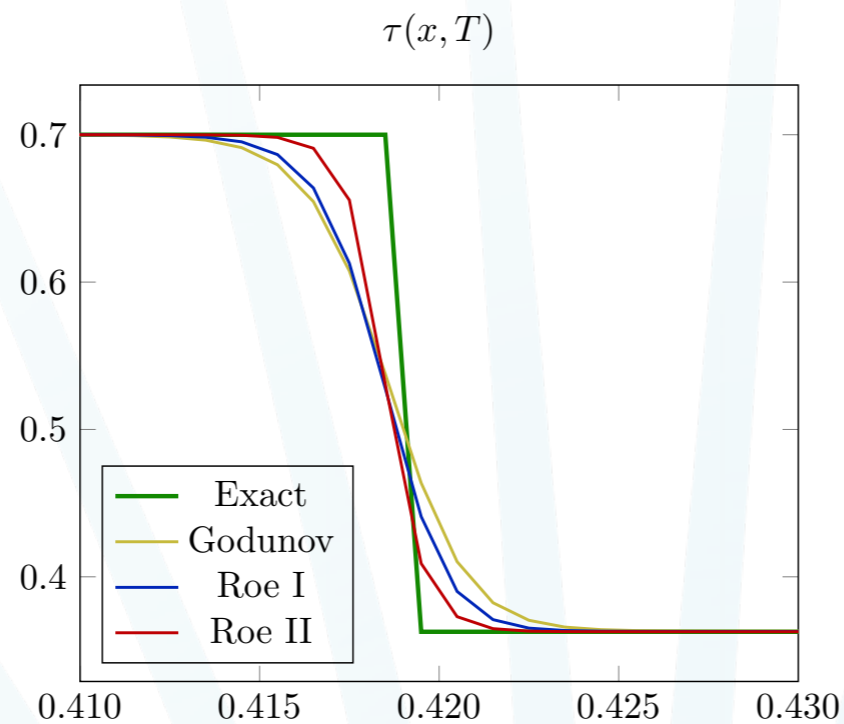
# Numerical results

Since the resolution of the state is a classical problem, all the schemes give good results.



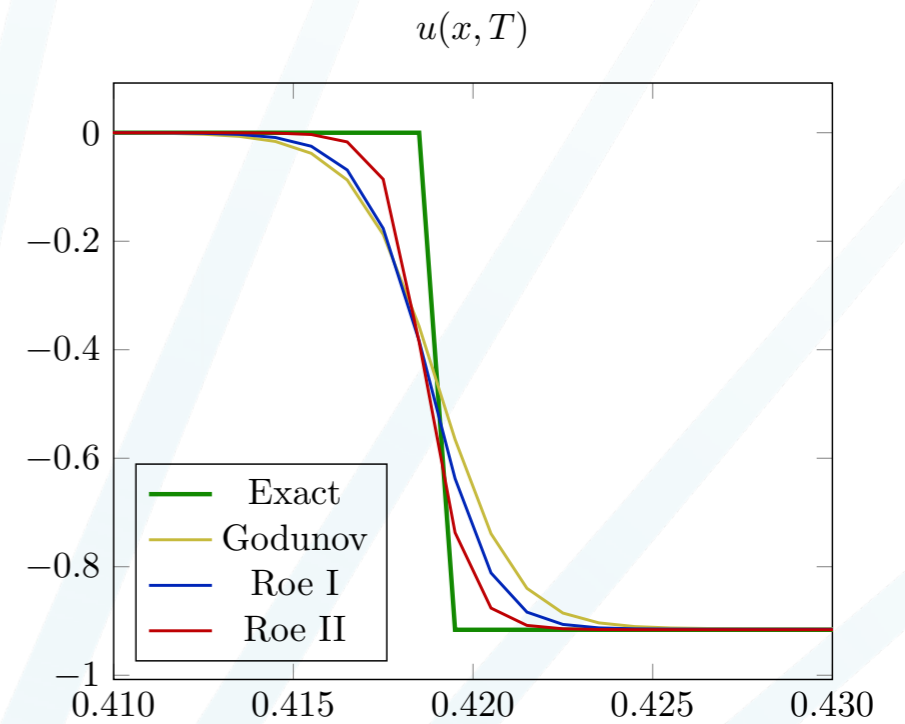
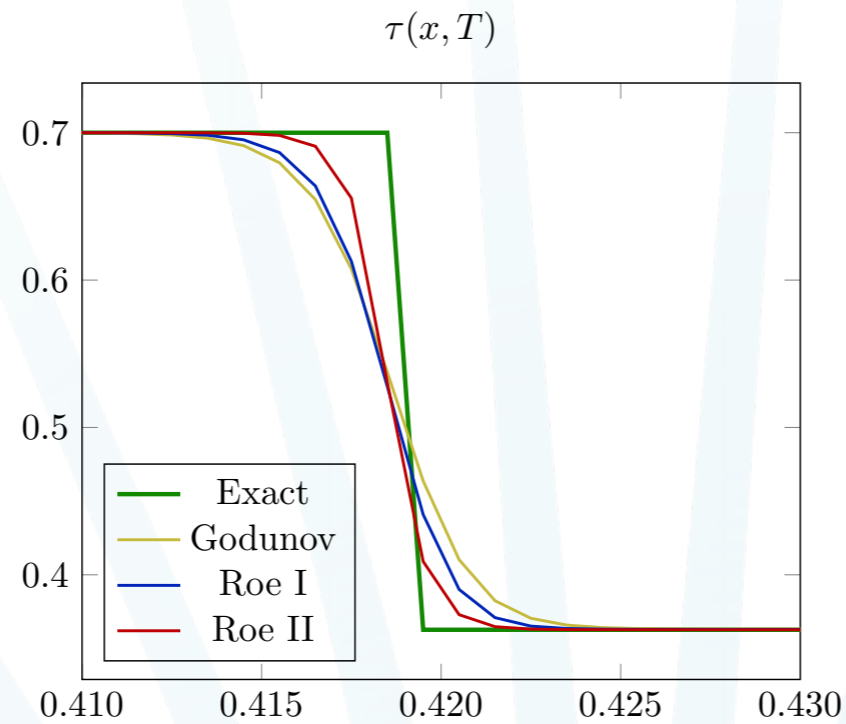
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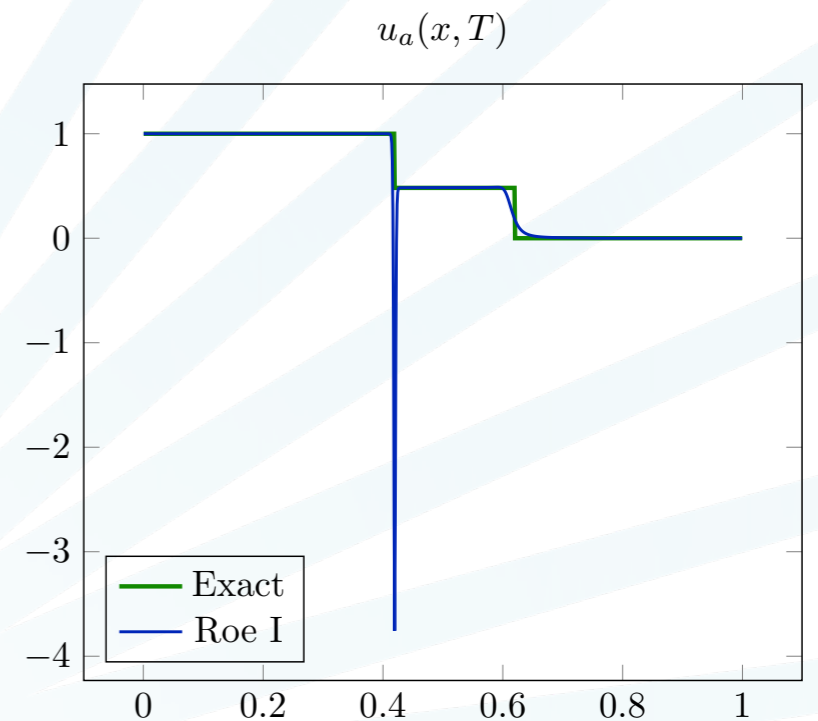
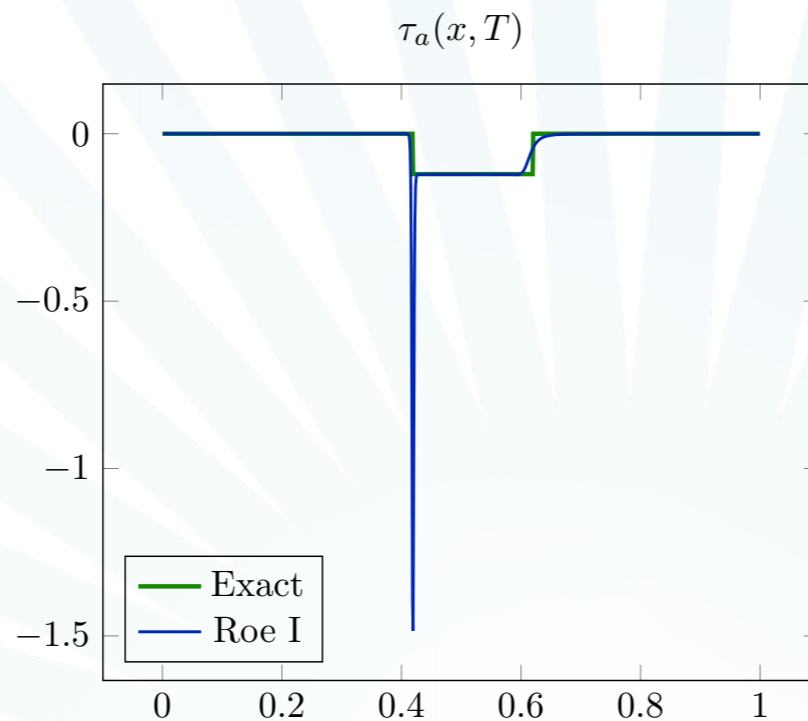


# Numerical results

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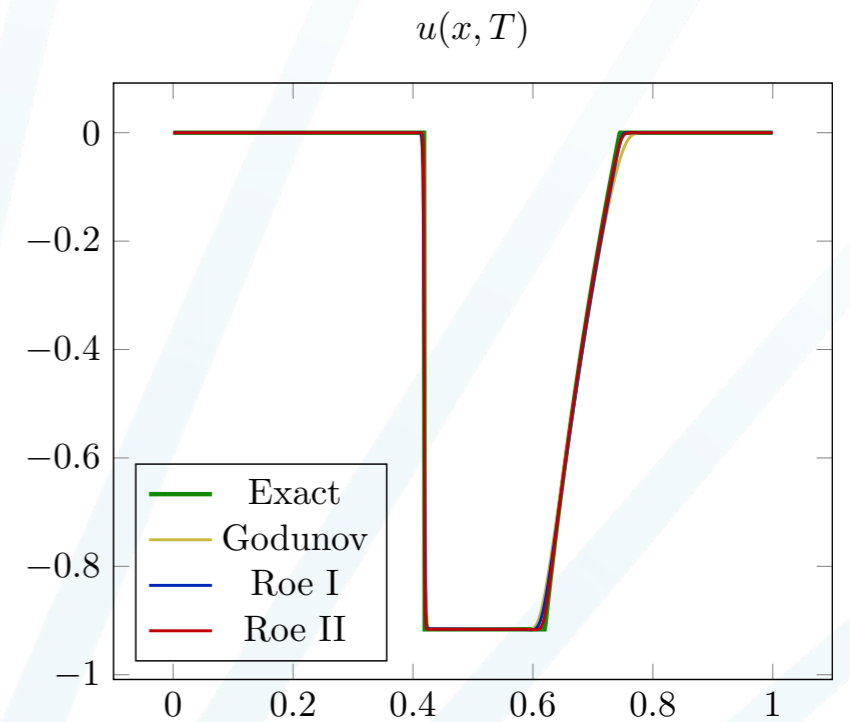
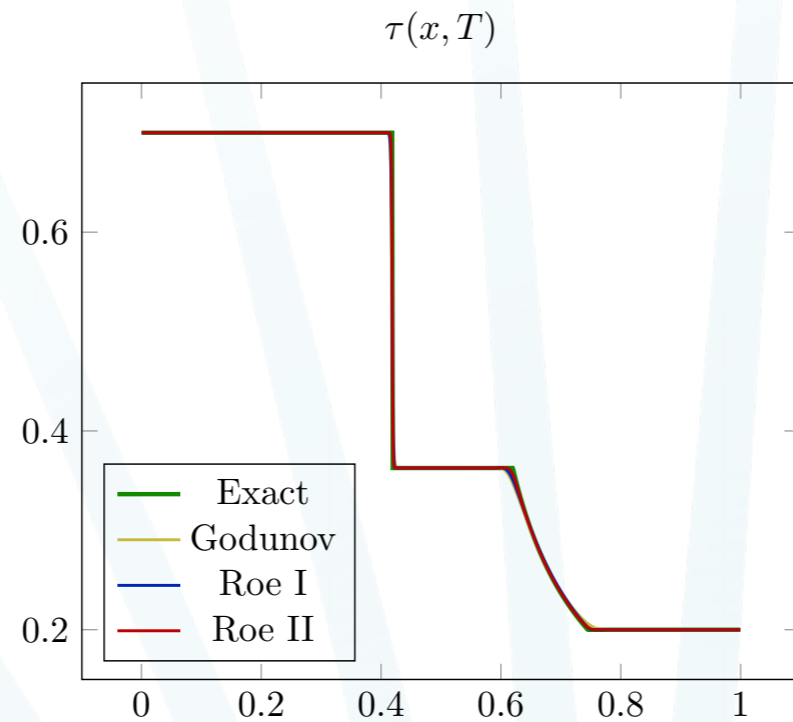


The same schemes do not work **without source term** for the sensitivity.

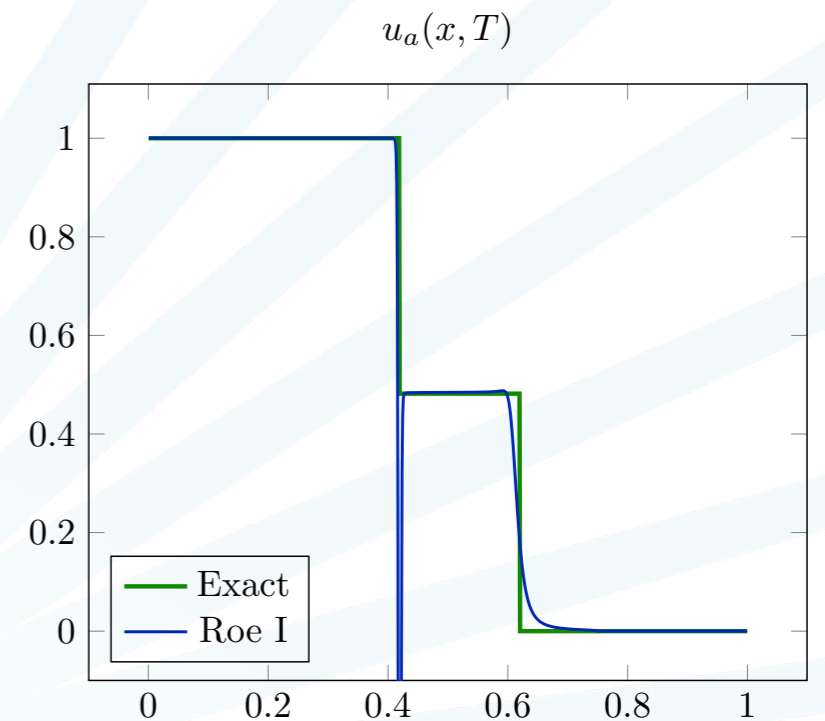
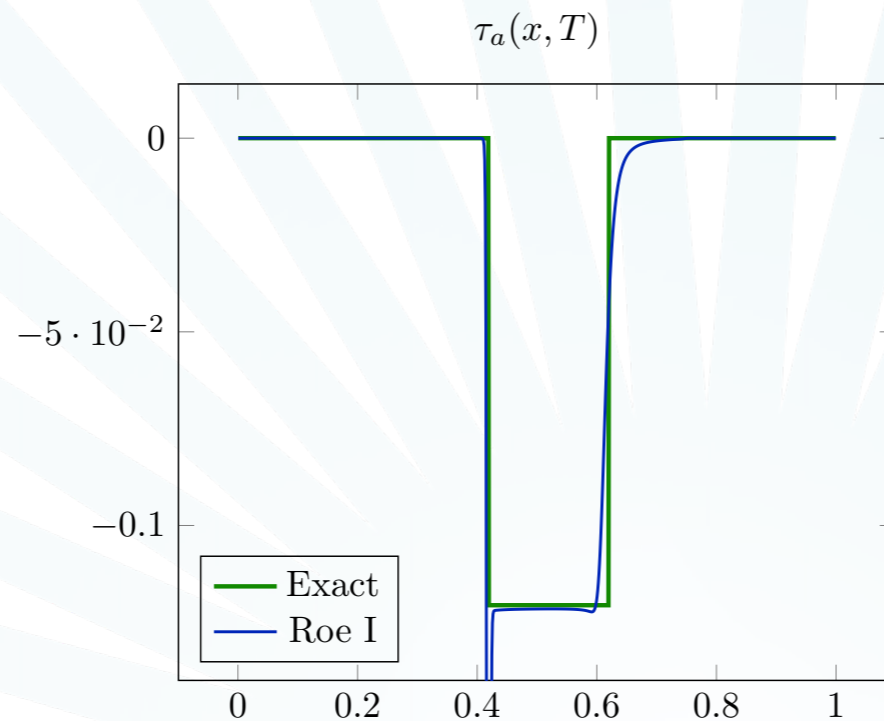


# Numerical results

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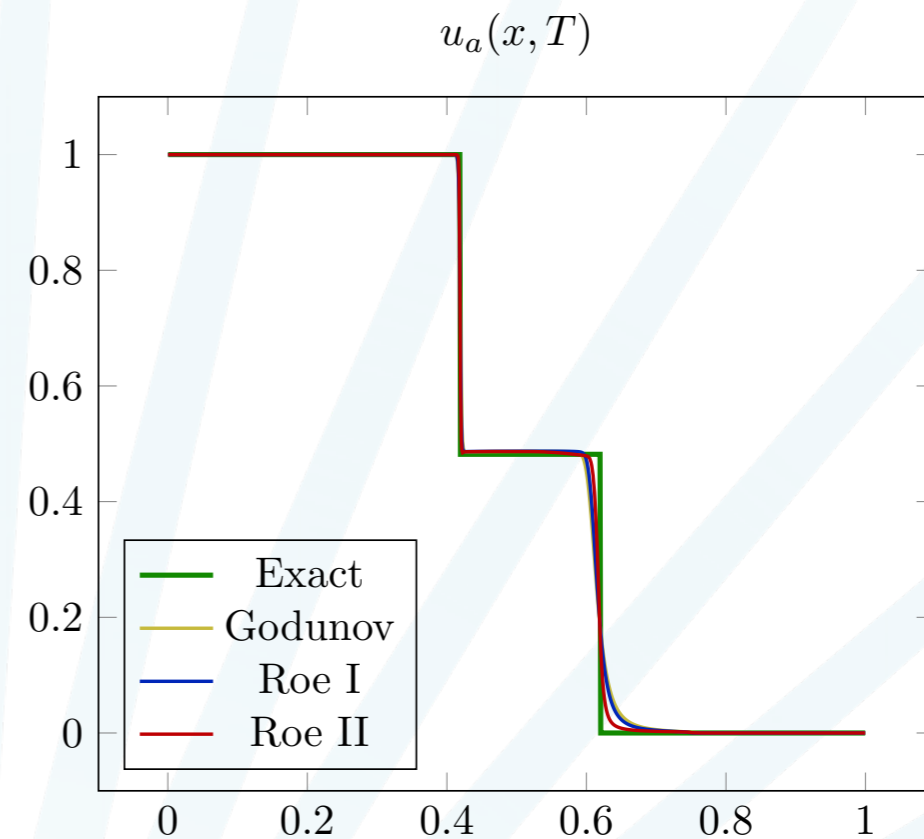
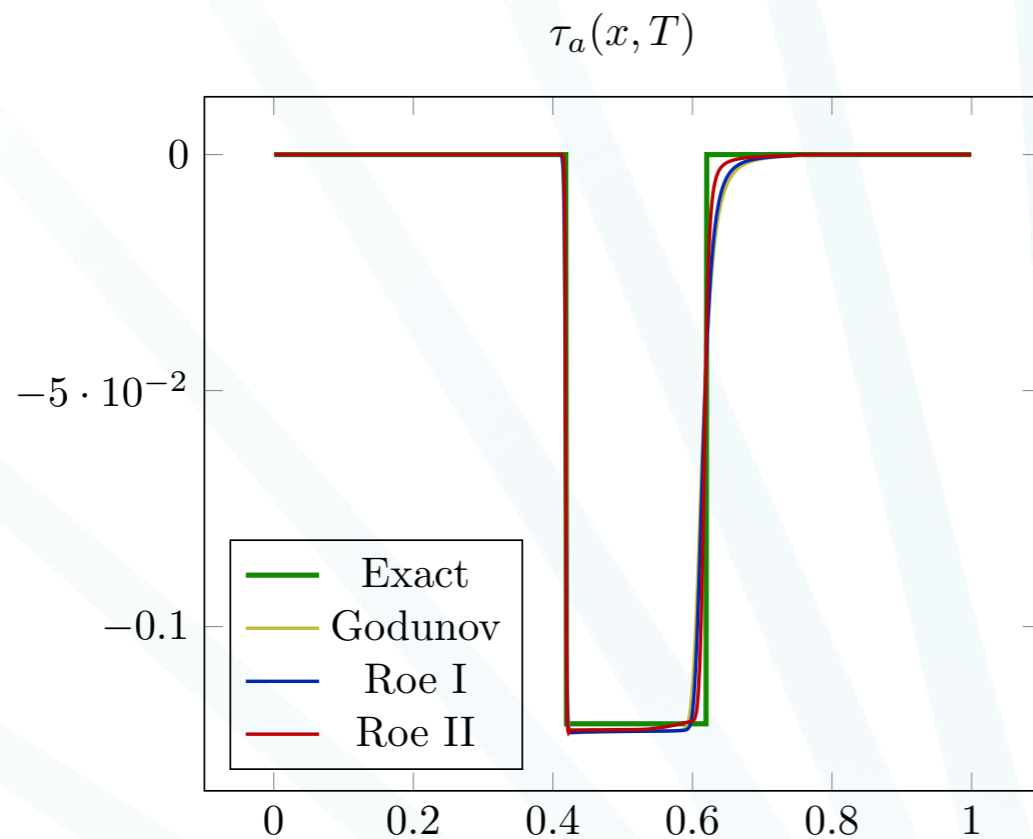


The same schemes do not work **without source term** for the sensitivity.



# Numerical results

The same schemes **with source term** for the sensitivity:

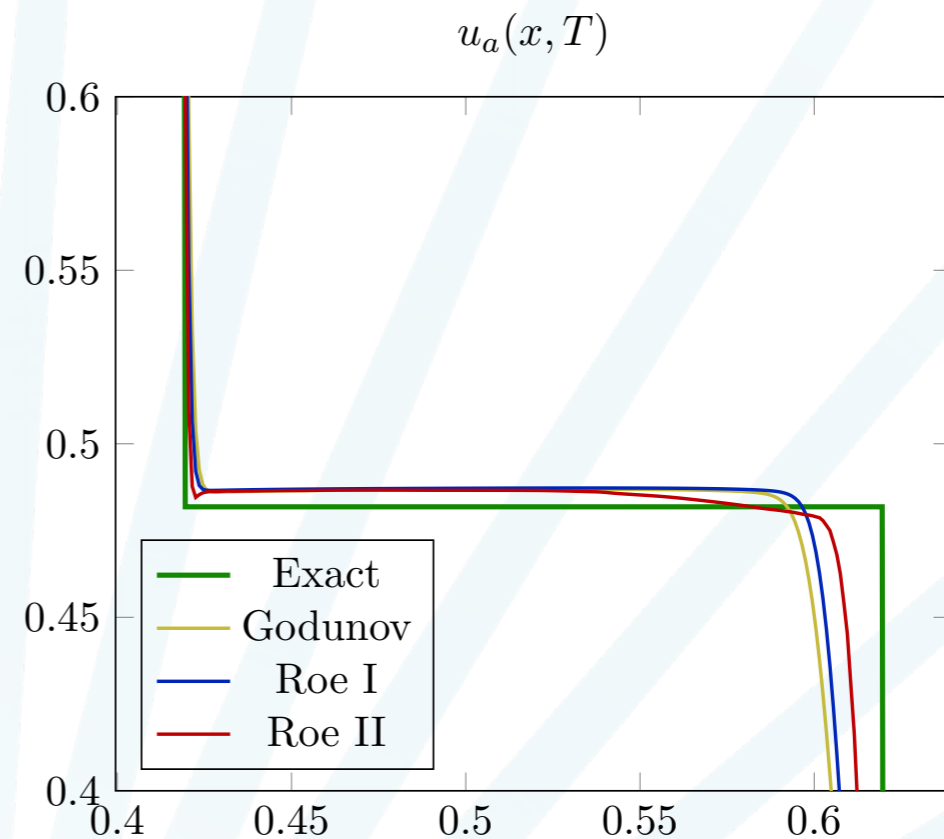
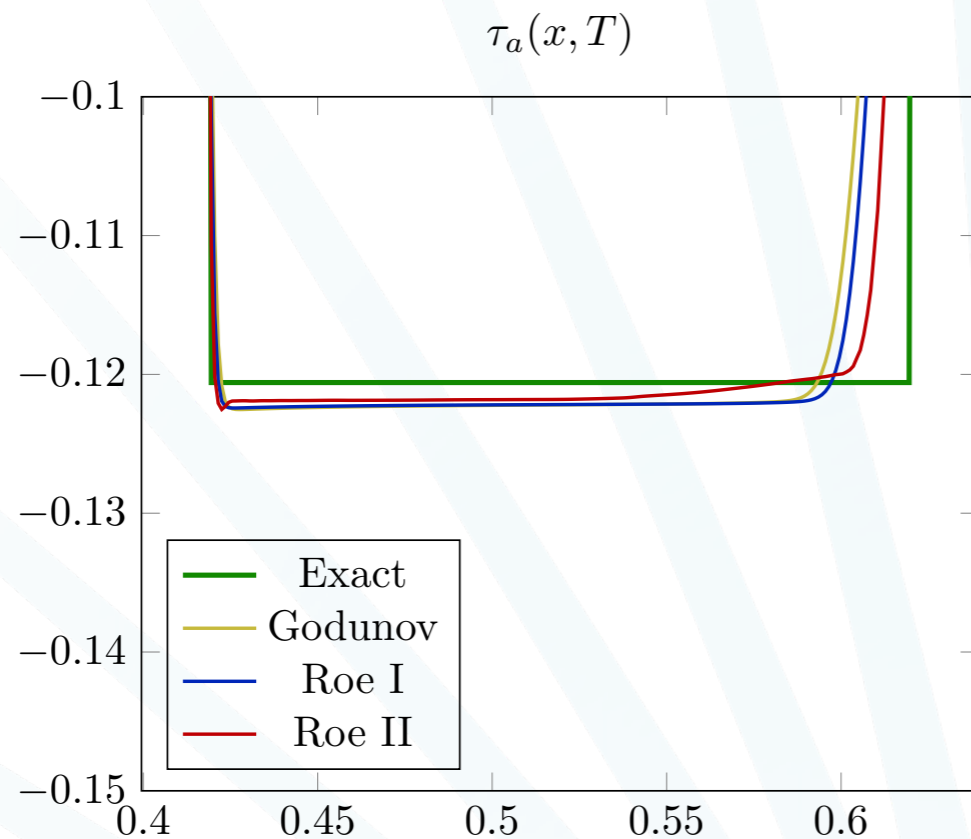


- Problems:
- ▶ the rarefaction is a discontinuity for the sensitivity,
  - ▶ the sensitivity value in the star zone is not correct.



# Numerical results

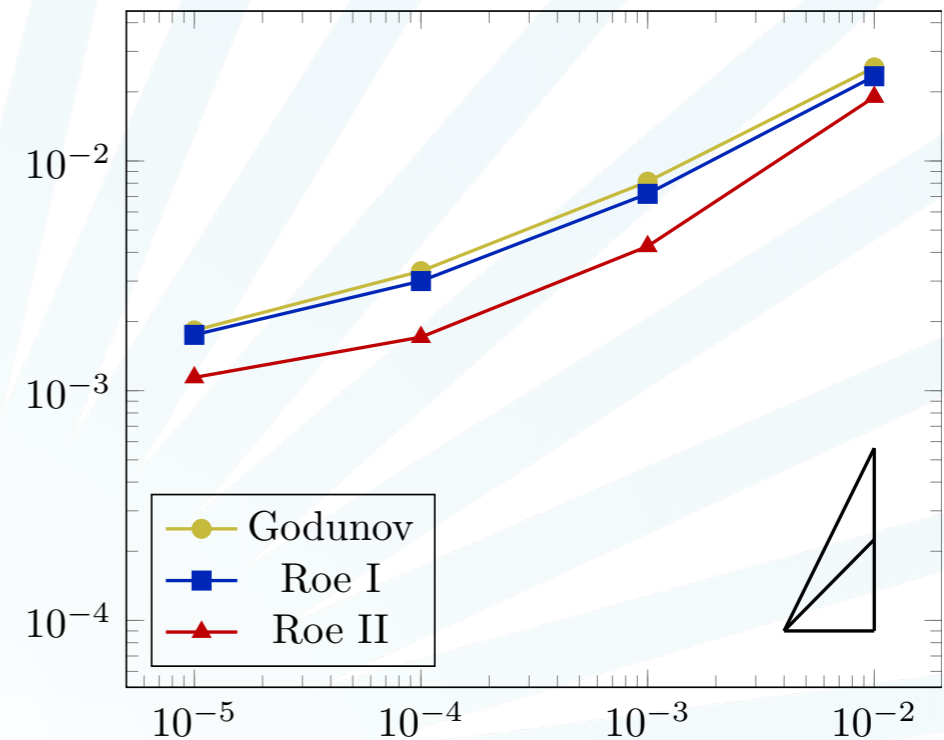
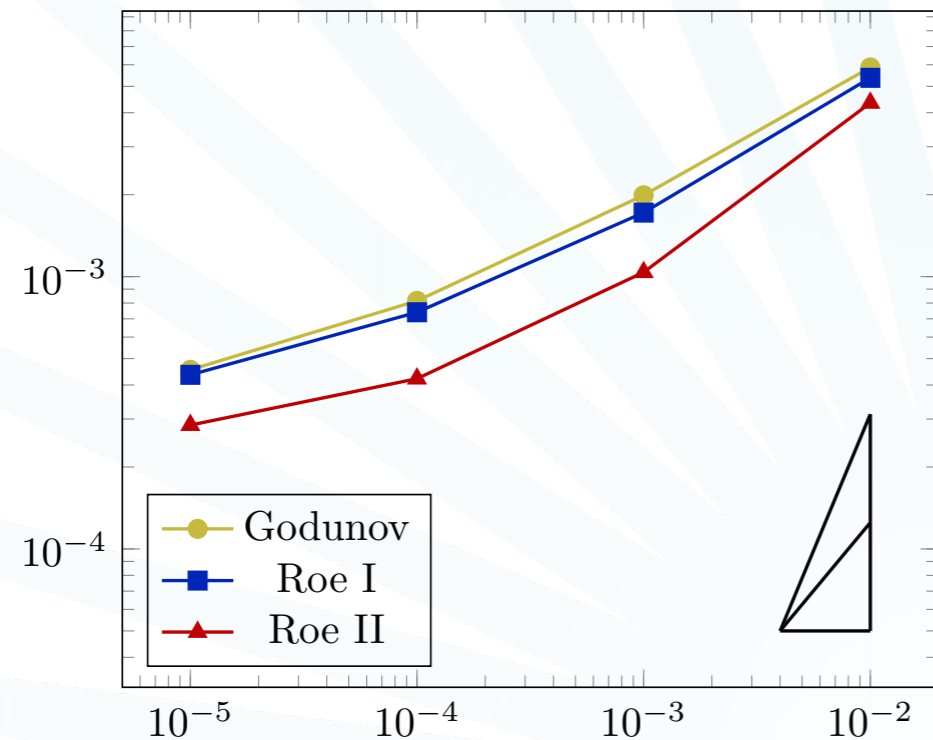
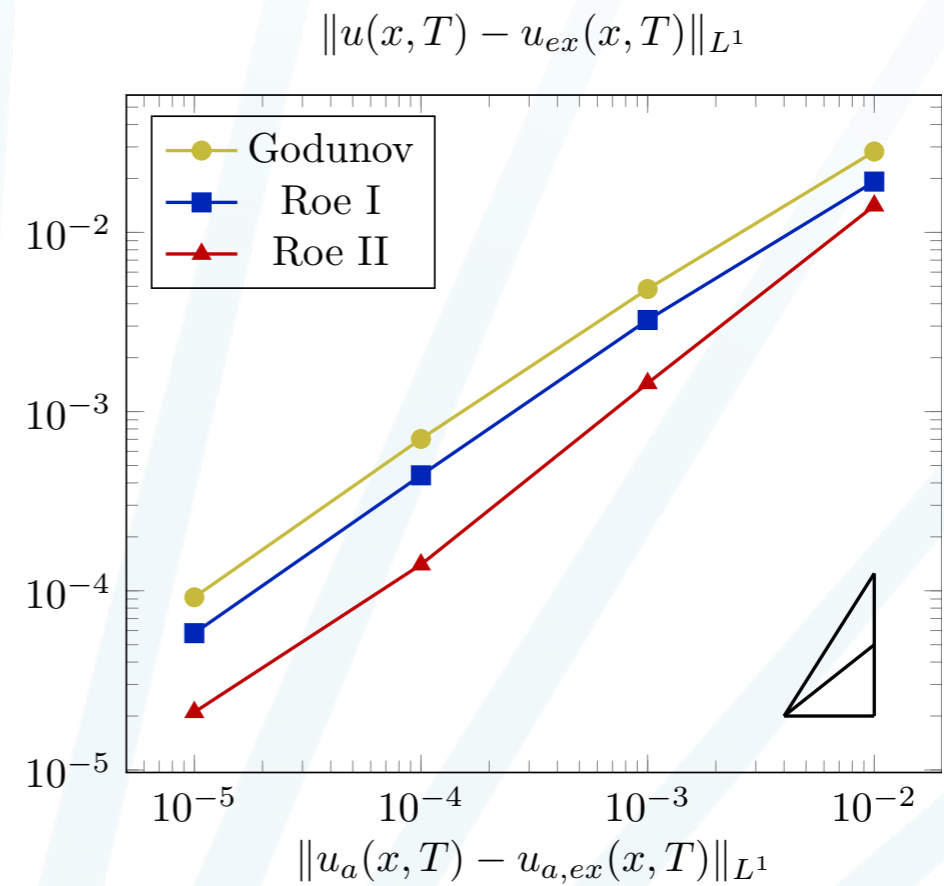
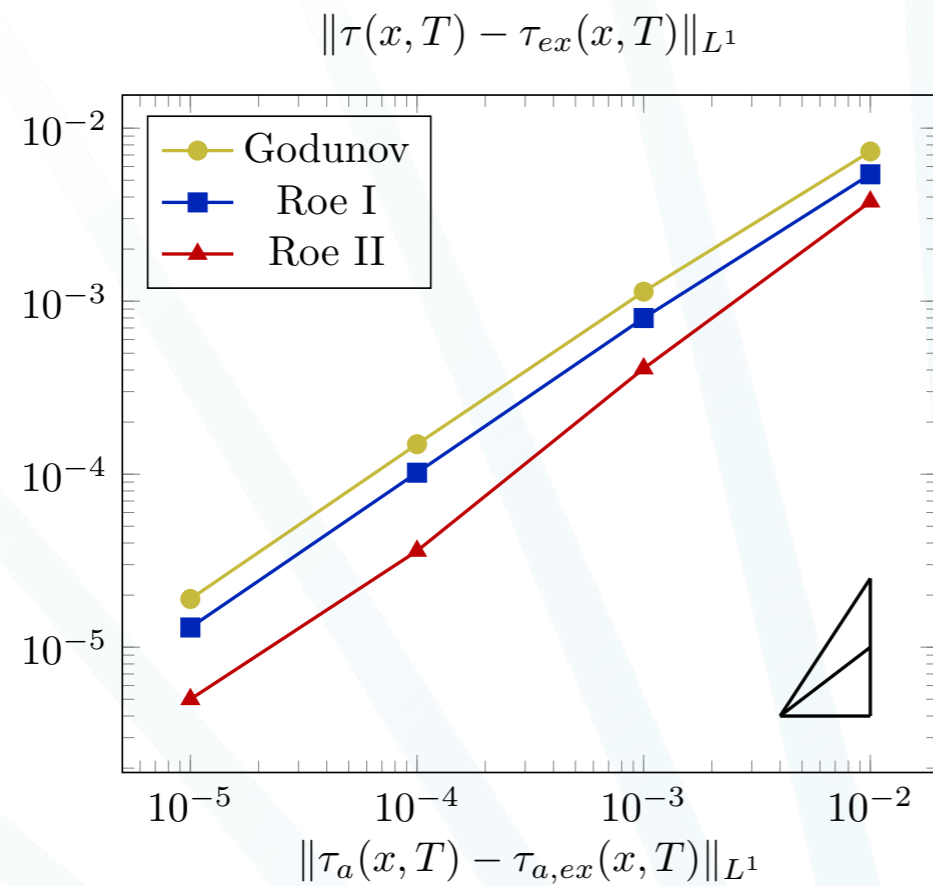
The same schemes **with source term** for the sensitivity:

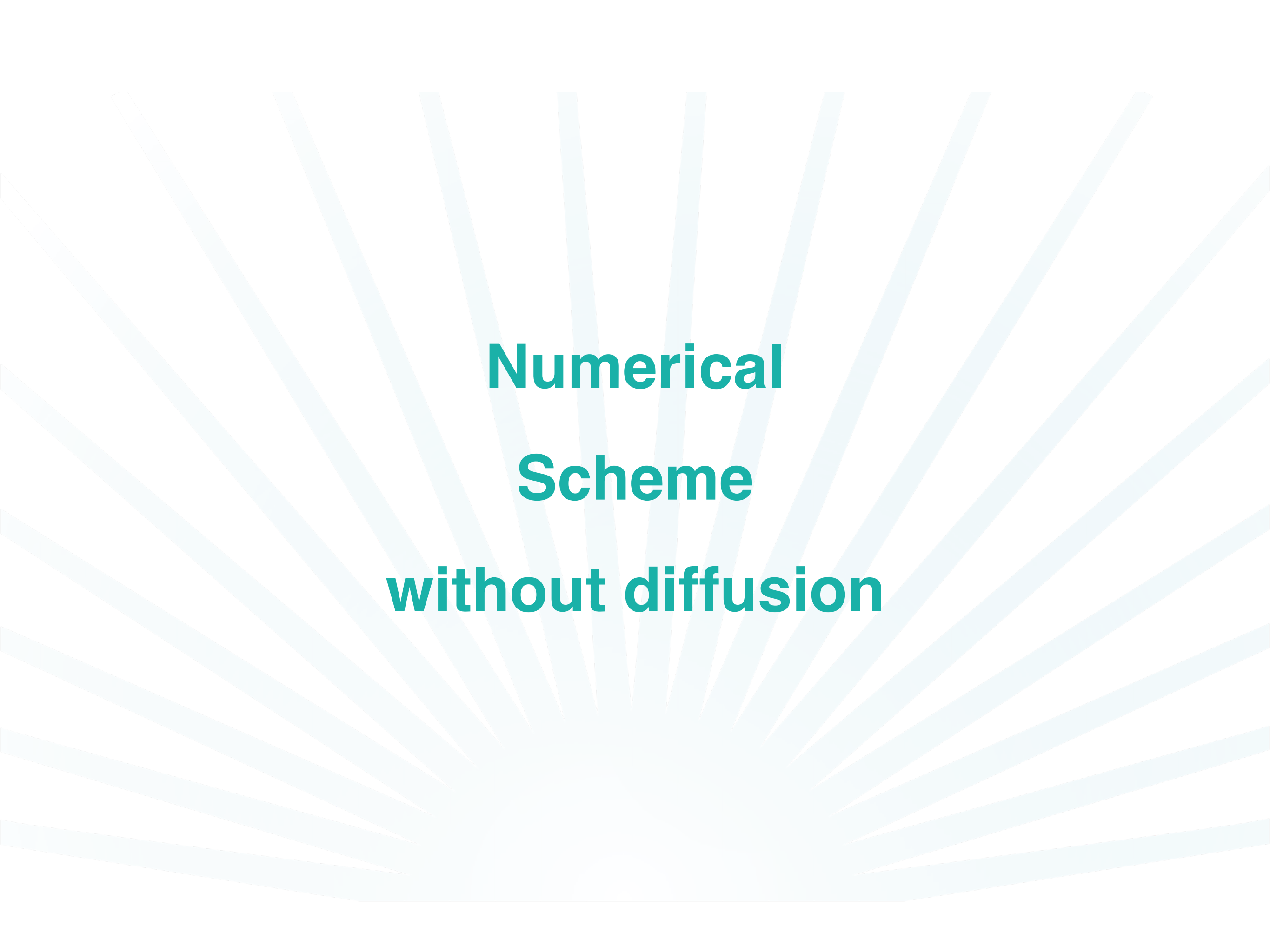


- Problems:
- ▶ the rarefaction is a discontinuity for the sensitivity,
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# Convergence

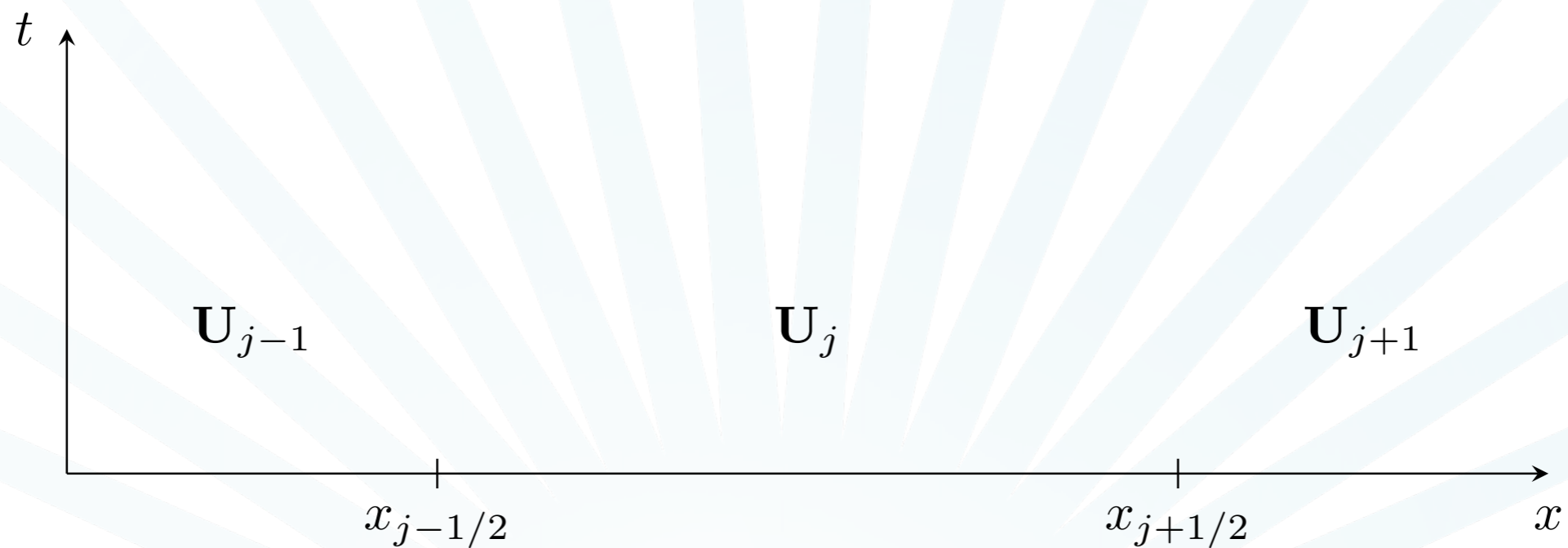




**Numerical  
Scheme  
without diffusion**

# Scheme without numerical diffusion

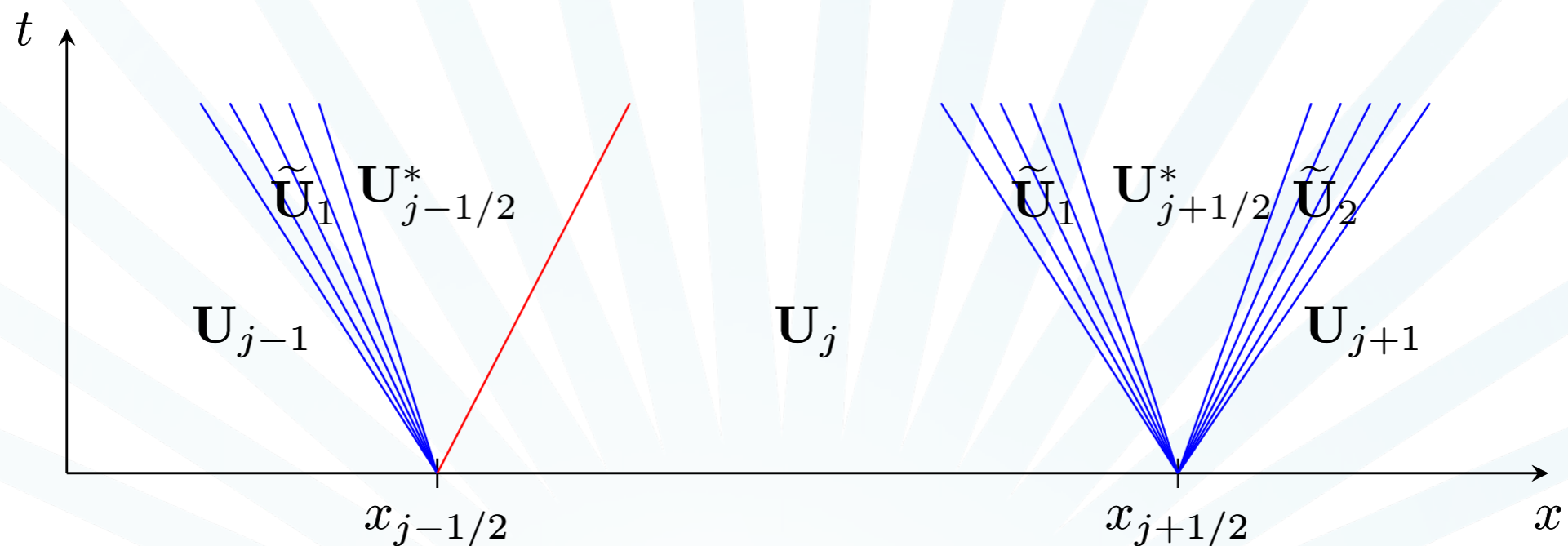
Step 0 : initial data discretisation



# Scheme without numerical diffusion

Step 0 : initial data discretisation

Step 1 : solution of the Riemann problems, one for each interface

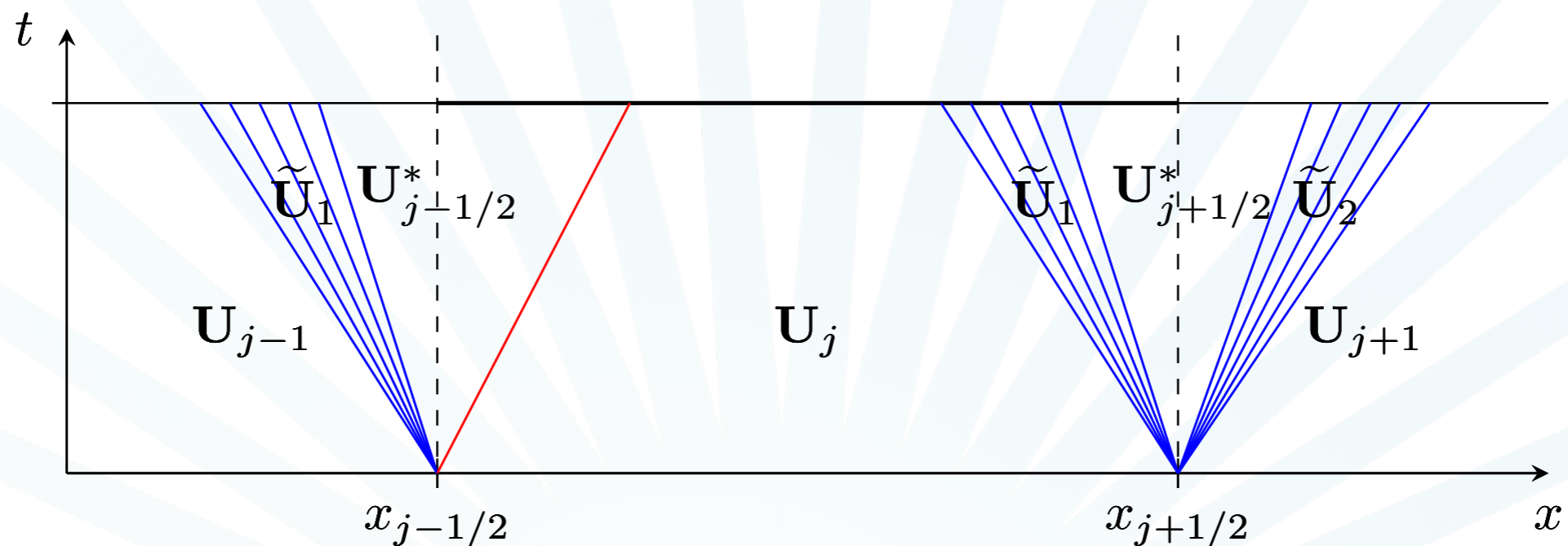


# Scheme without numerical diffusion

Step 0 : initial data discretisation

Step 1 : solution of the Riemann problems, one for each interface

Step 2 : average

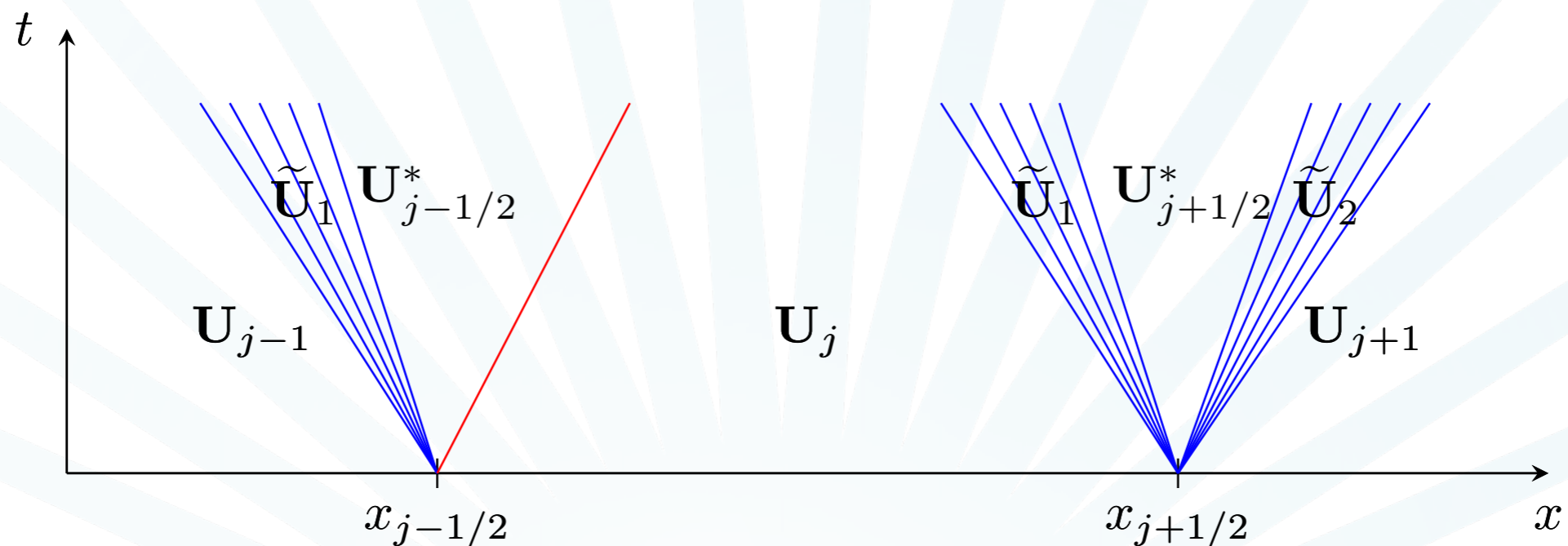


# Scheme without numerical diffusion

Step 0 : initial data discretisation

Step 1 : solution of the Riemann problems, one for each interface

~~Step 2 : average~~



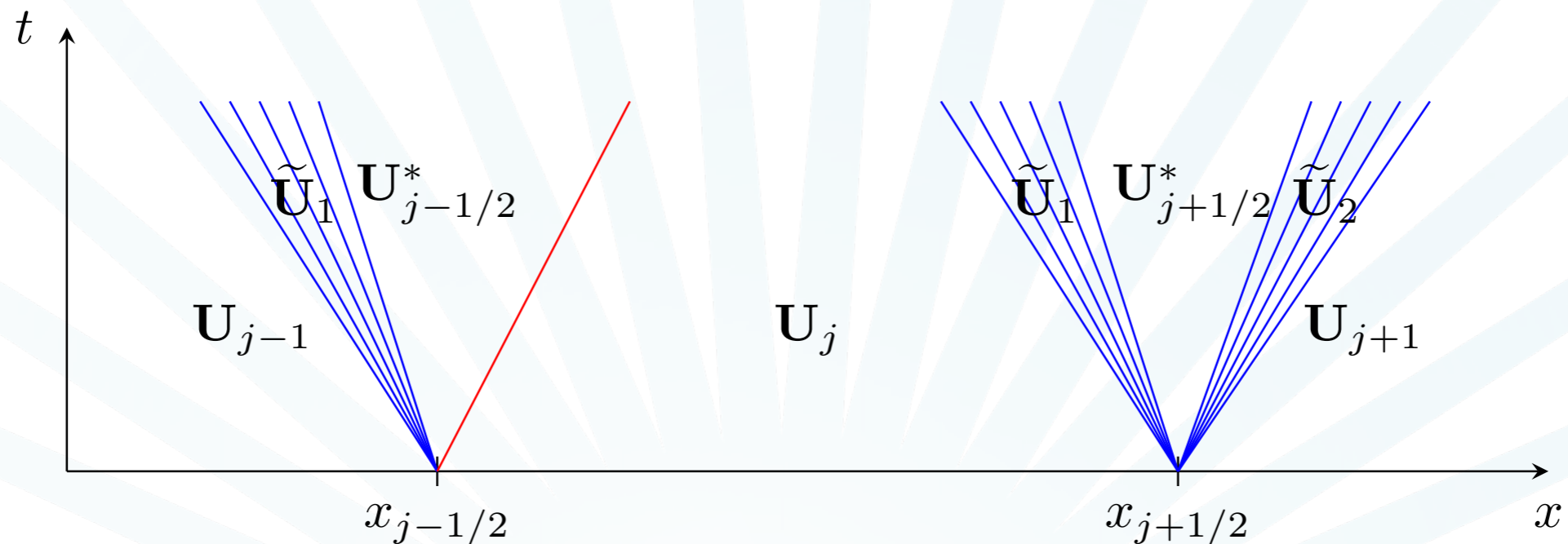


# Scheme without numerical diffusion

Step 0 : initial data discretisation

Step 1 : solution of the Riemann problems, one for each interface

Step 2 : definition of a staggered mesh on which the average is performed

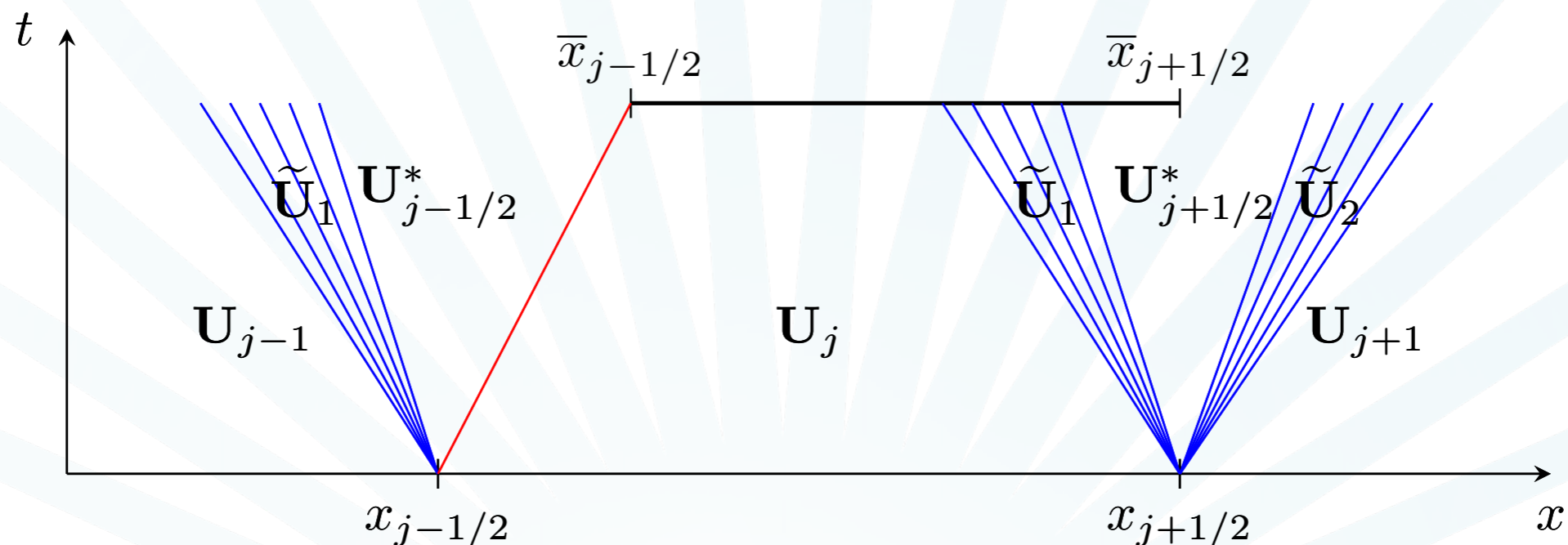


# Scheme without numerical diffusion

Step 0 : initial data discretisation

Step 1 : solution of the Riemann problems, one for each interface

Step 2 : definition of a staggered mesh on which the average is performed



$$\bar{x}_{j-1/2} = x_{j-1/2} + \sigma_{j-1/2} \Delta t$$

# Scheme without numerical diffusion

Step 0 : initial data discretisation

Step 1 : solution of the Riemann problems, one for each interface

Step 2 : definition of a staggered mesh on which the average is performed

Step 3 : projection on the initial mesh

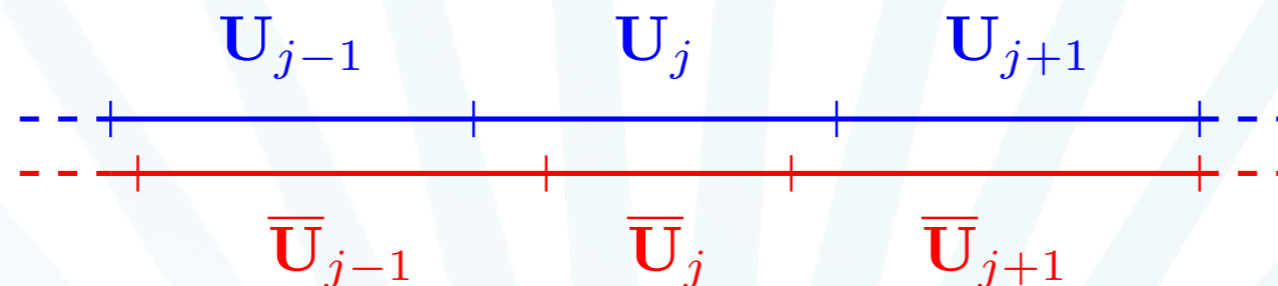
# Scheme without numerical diffusion

Step 0 : initial data discretisation

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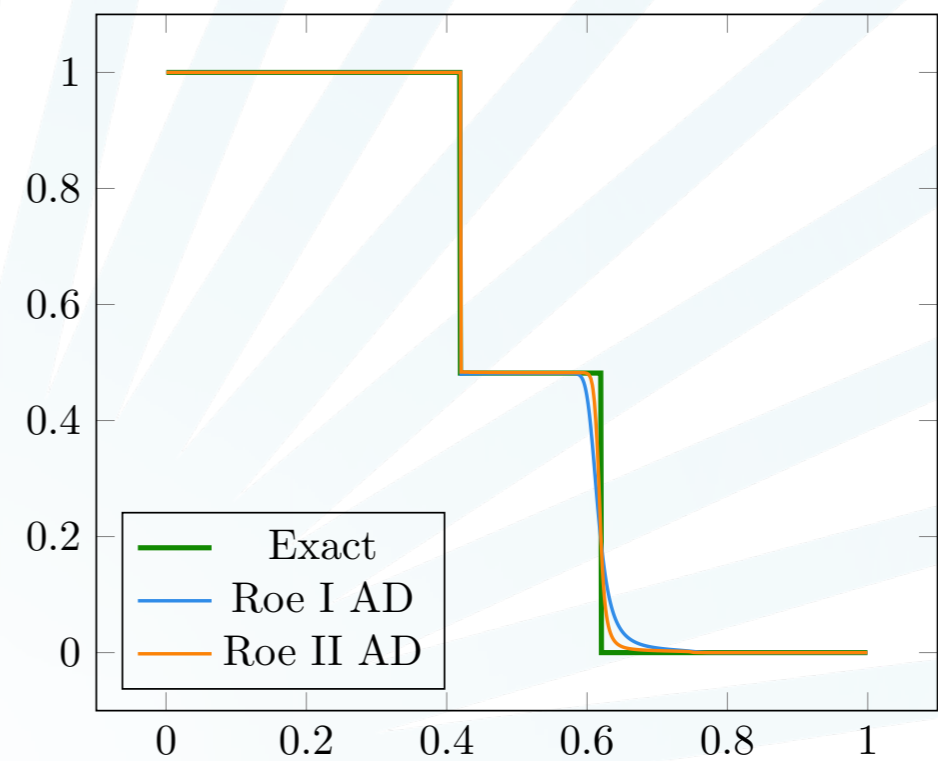
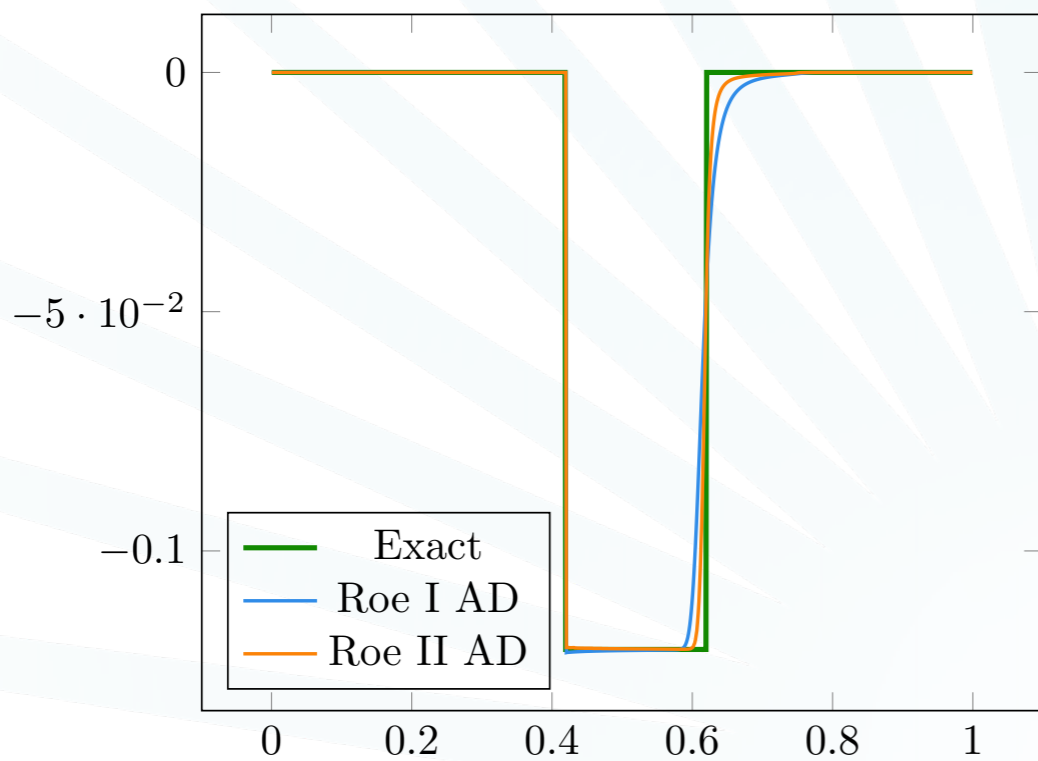
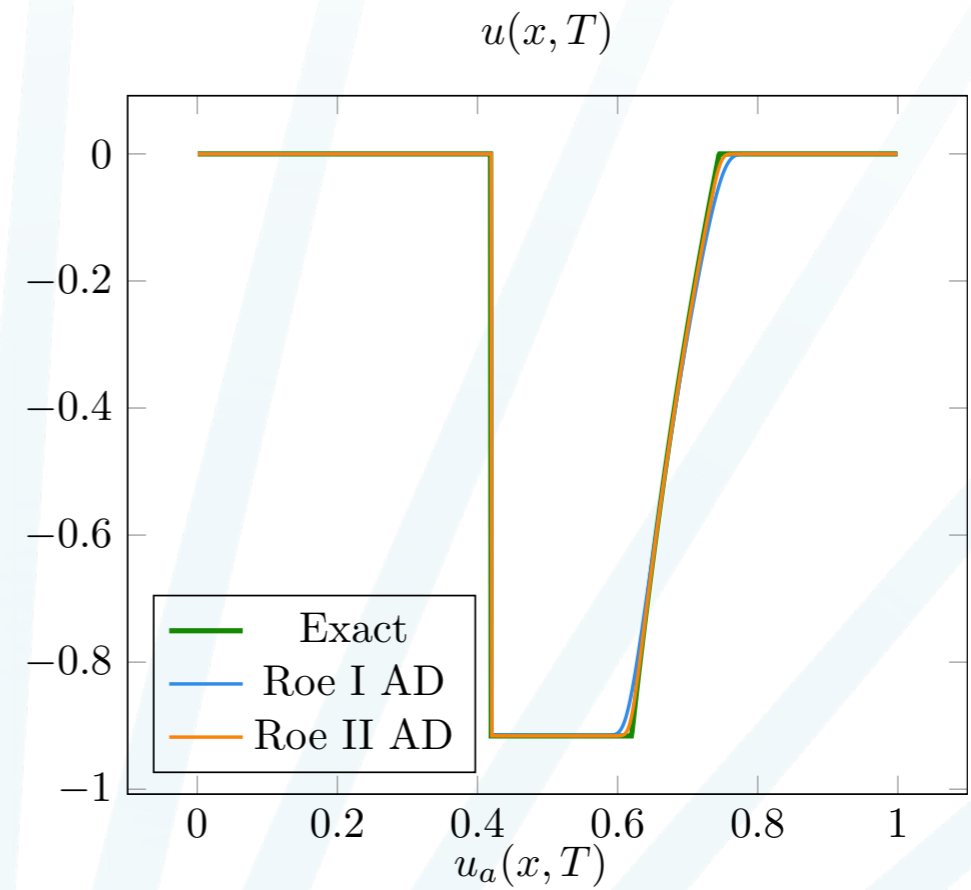
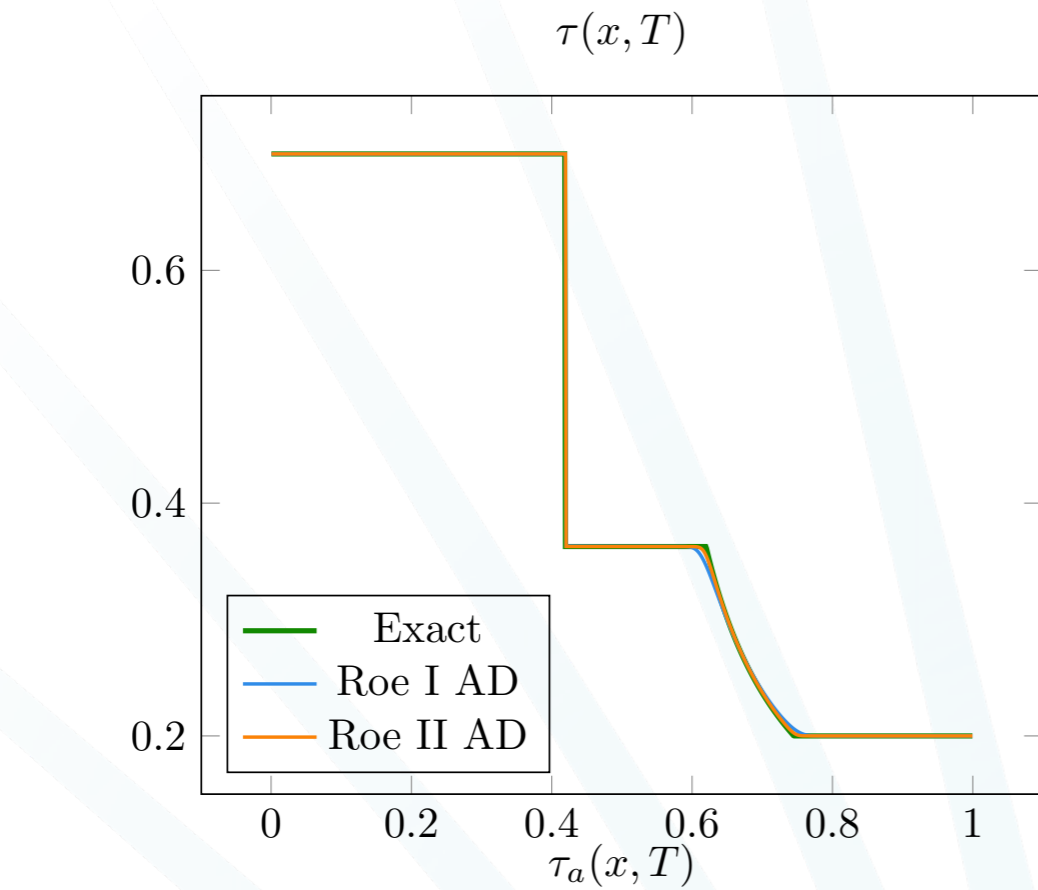
Step 3 : projection on the initial mesh



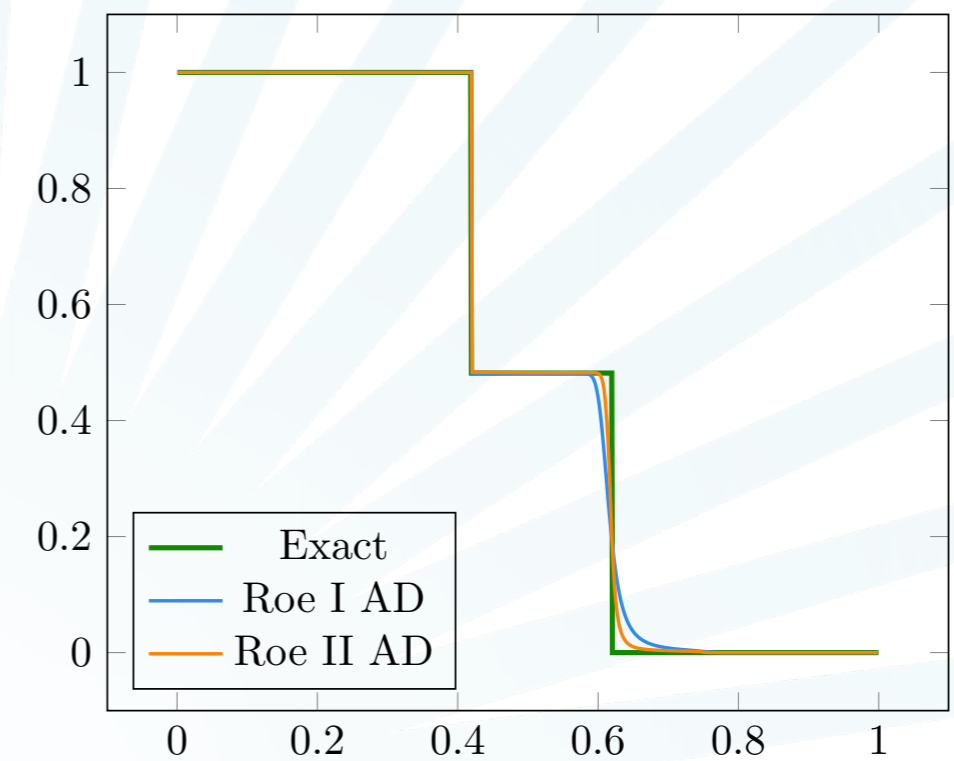
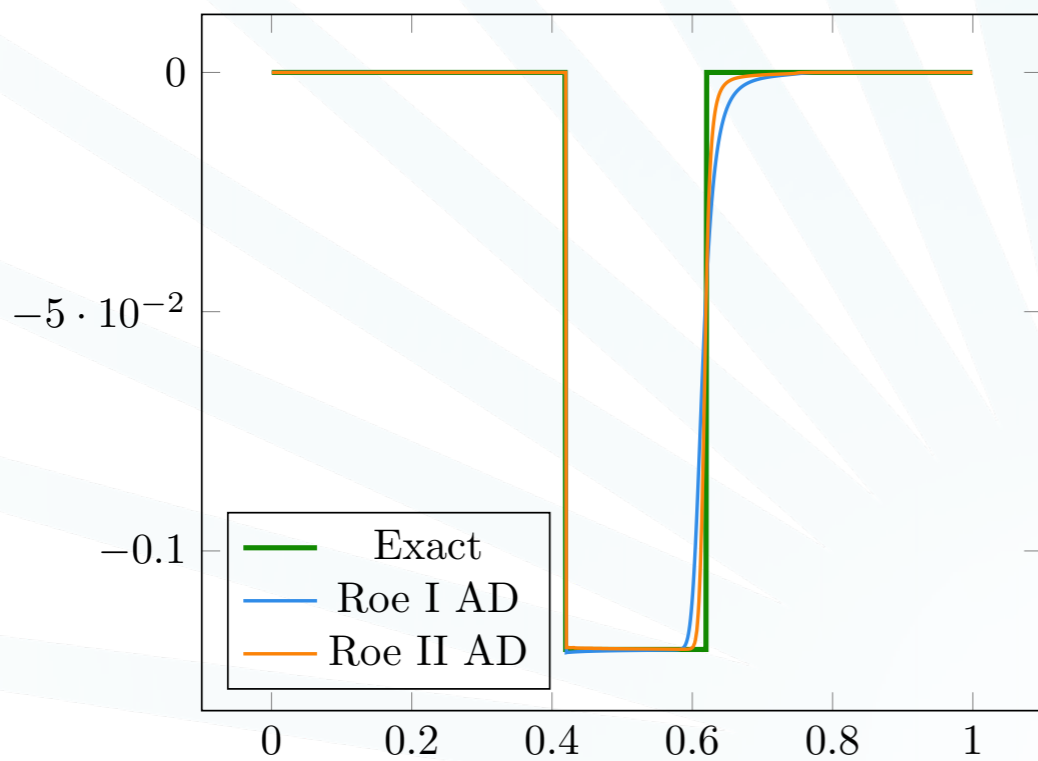
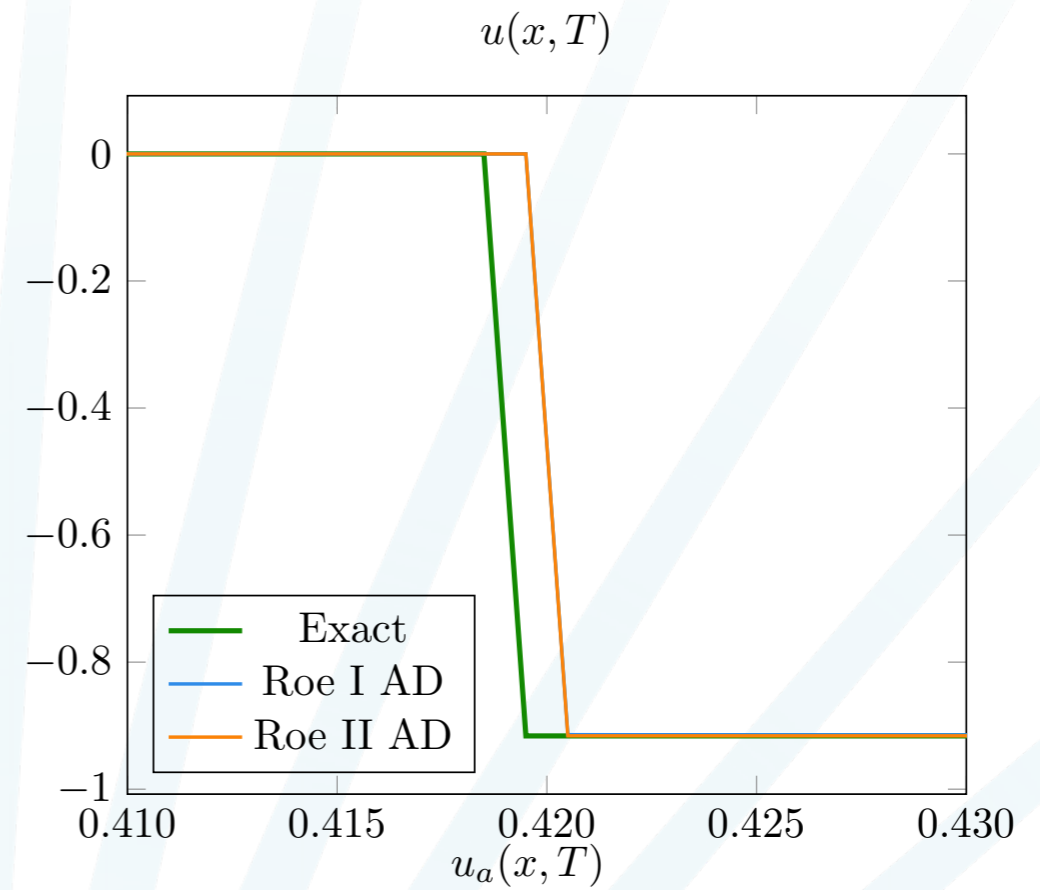
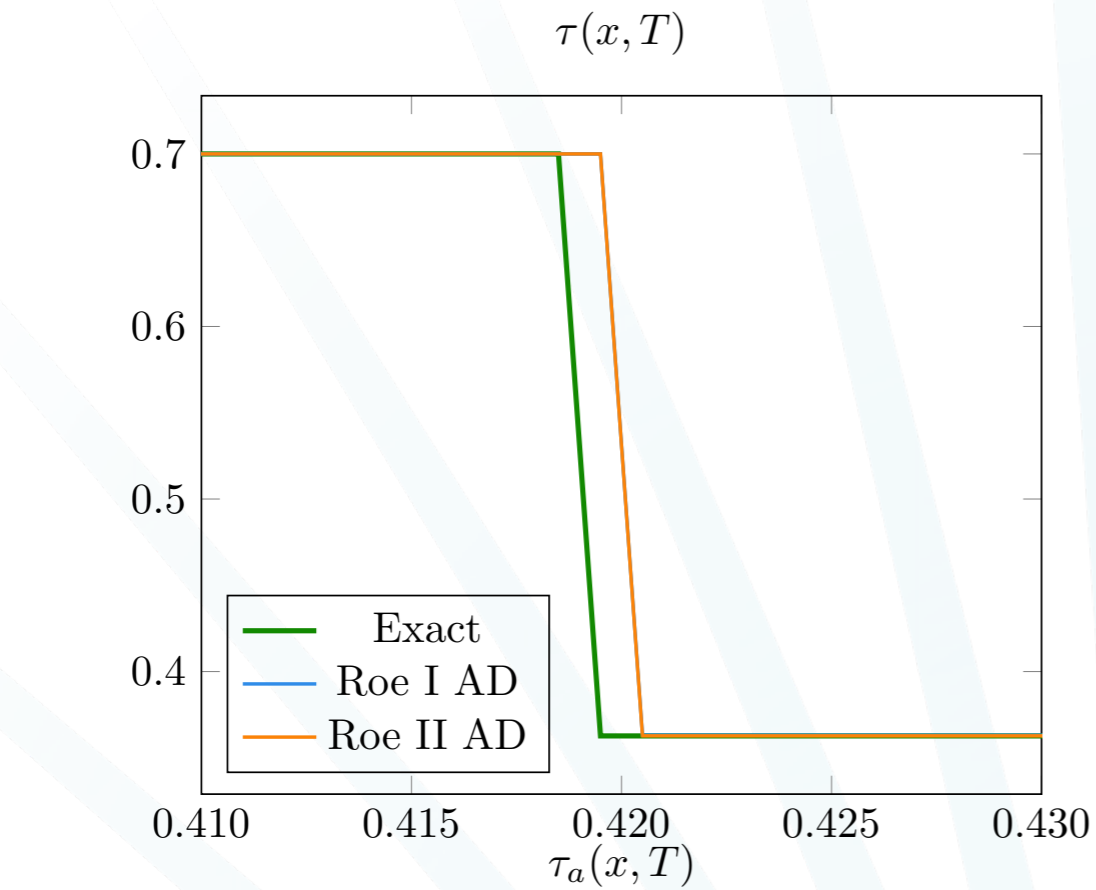
$$\mathbf{U}_j = \begin{cases} \bar{\mathbf{U}}_{j-1} & \text{if } a \in \left(0, \frac{\Delta t}{\Delta x} \max(\sigma_{j-1/2}, 0)\right), \\ \bar{\mathbf{U}}_j & \text{if } a \in \left[\frac{\Delta t}{\Delta x} \max(\sigma_{j-1/2}, 0), 1 + \frac{\Delta t}{\Delta x} \min(\sigma_{j+1/2}, 0)\right), \\ \bar{\mathbf{U}}_{j+1} & \text{if } a \in \left[1 + \frac{\Delta t}{\Delta x} \min(\sigma_{j+1/2}, 0), 1\right). \end{cases}$$

$$a \sim \mathcal{U}([0, 1])$$

# Results

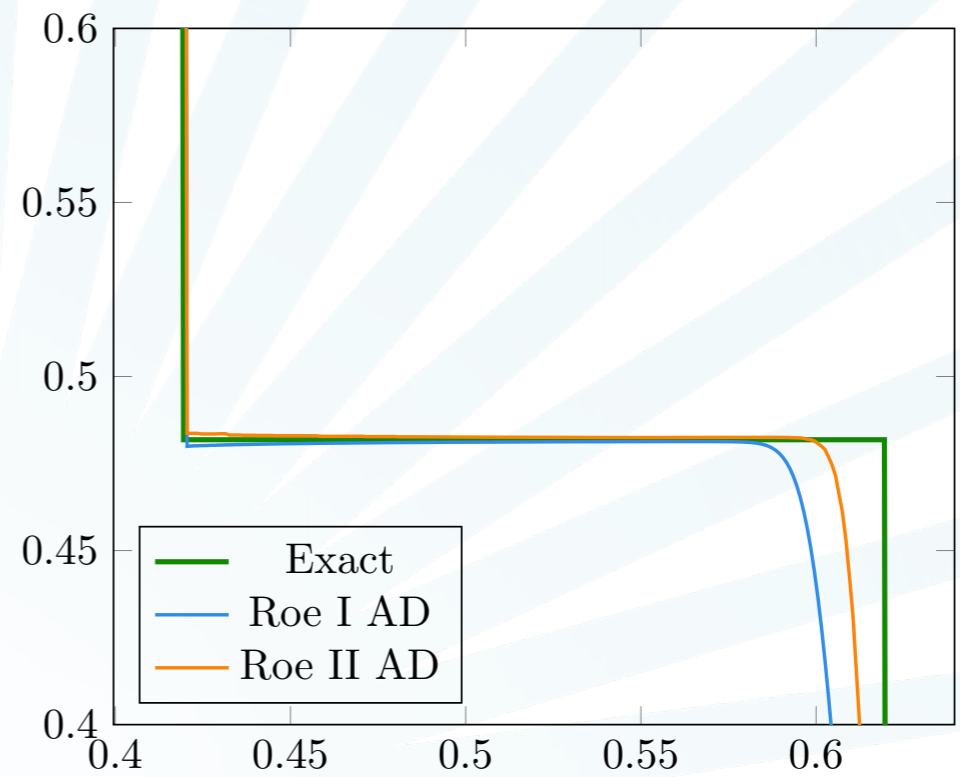
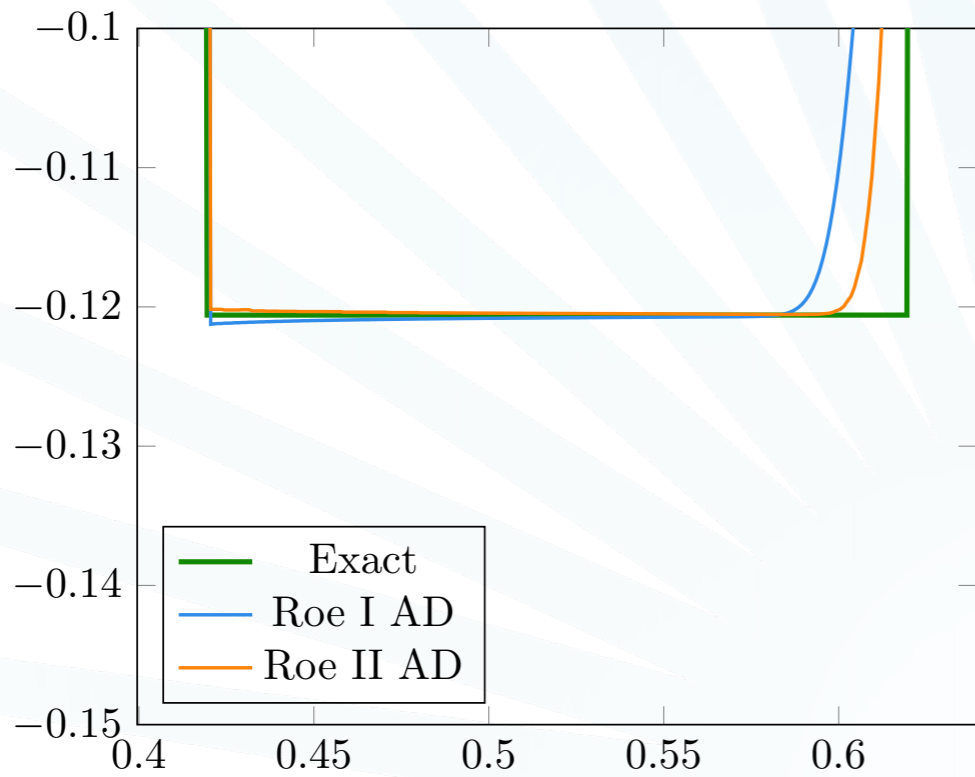
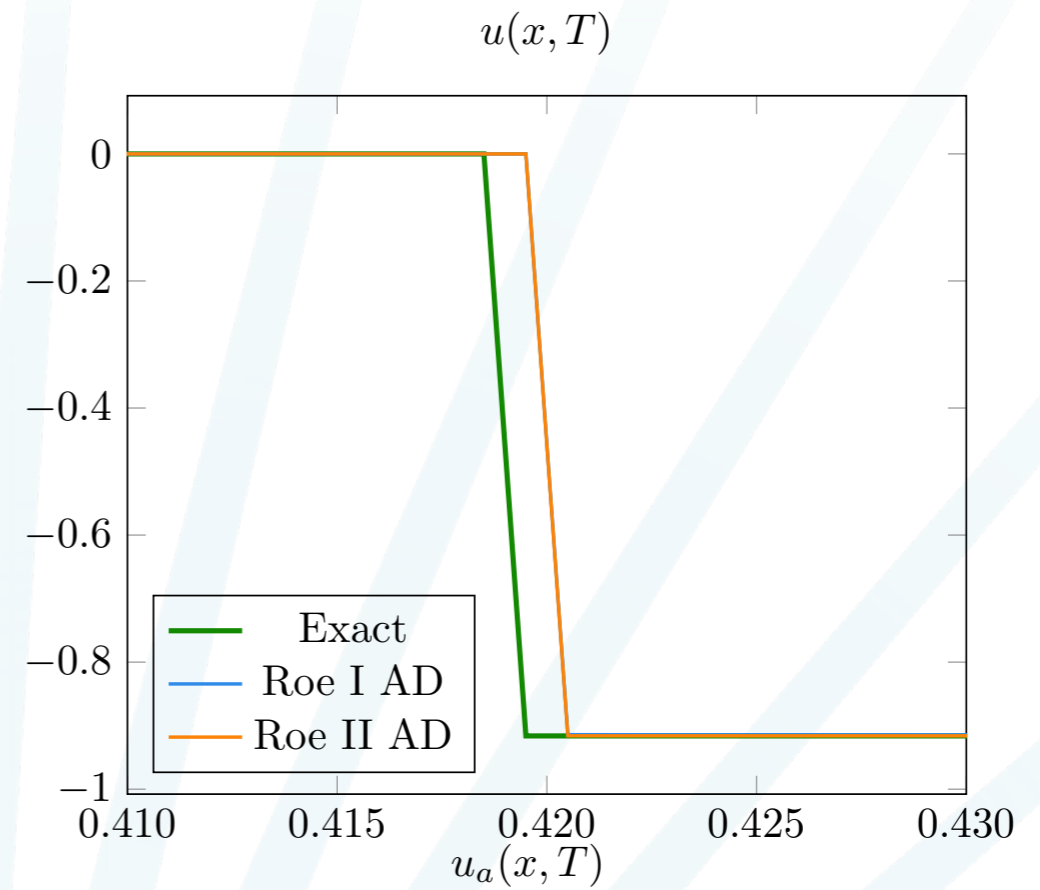
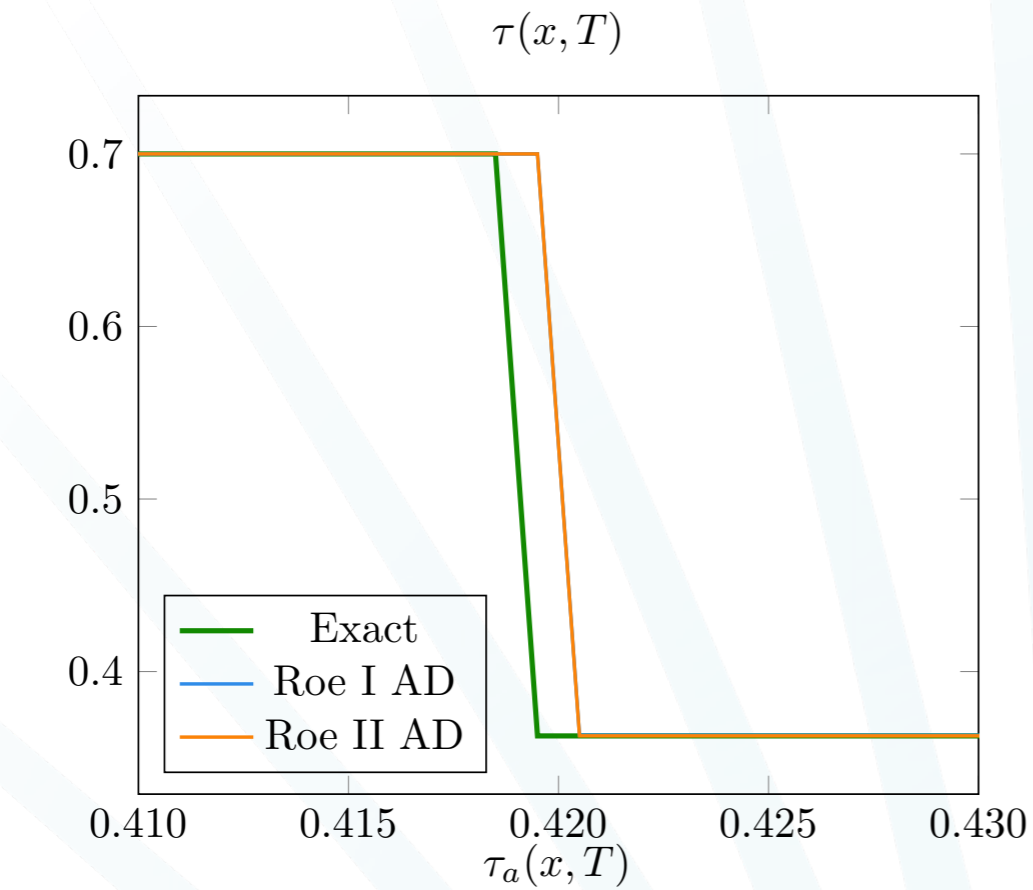


# Results

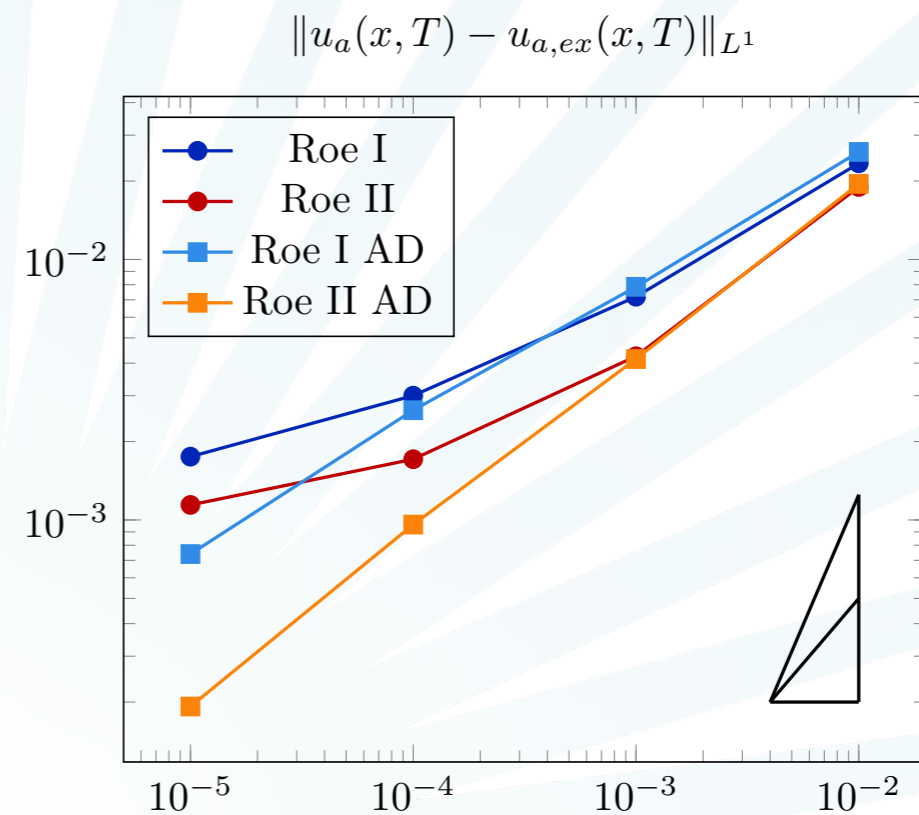
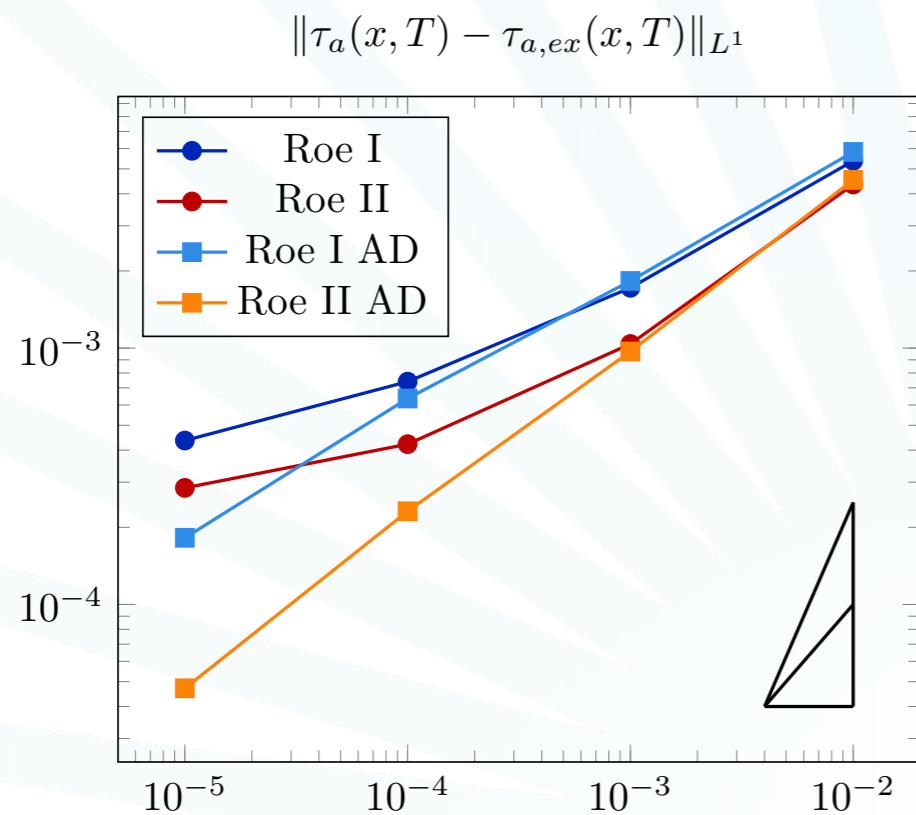
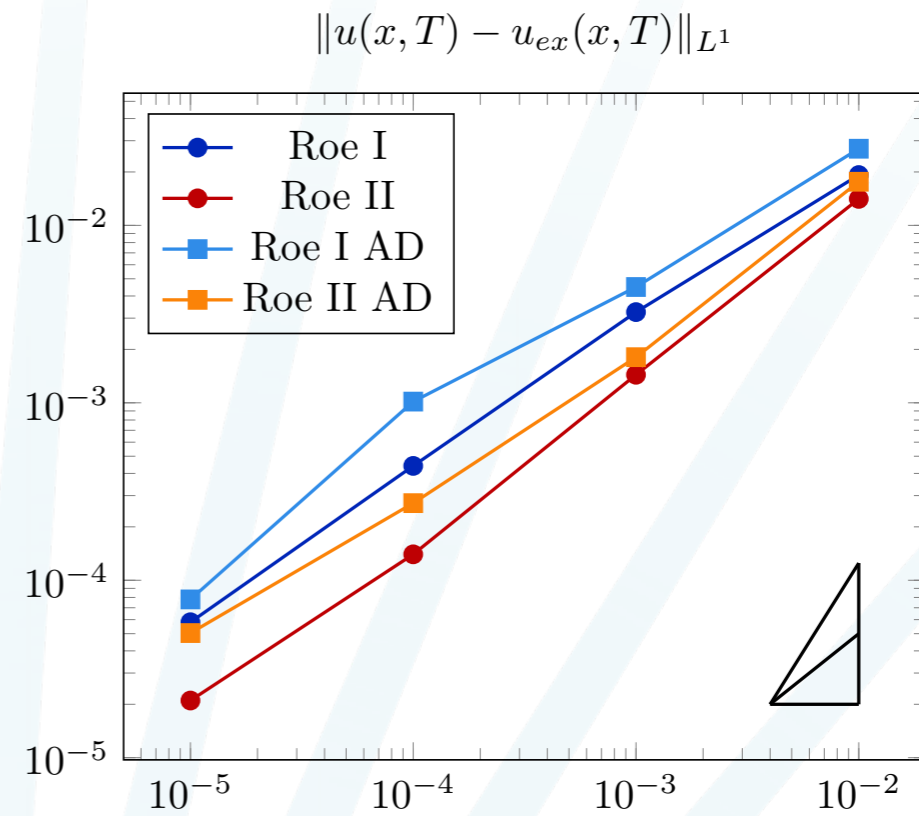
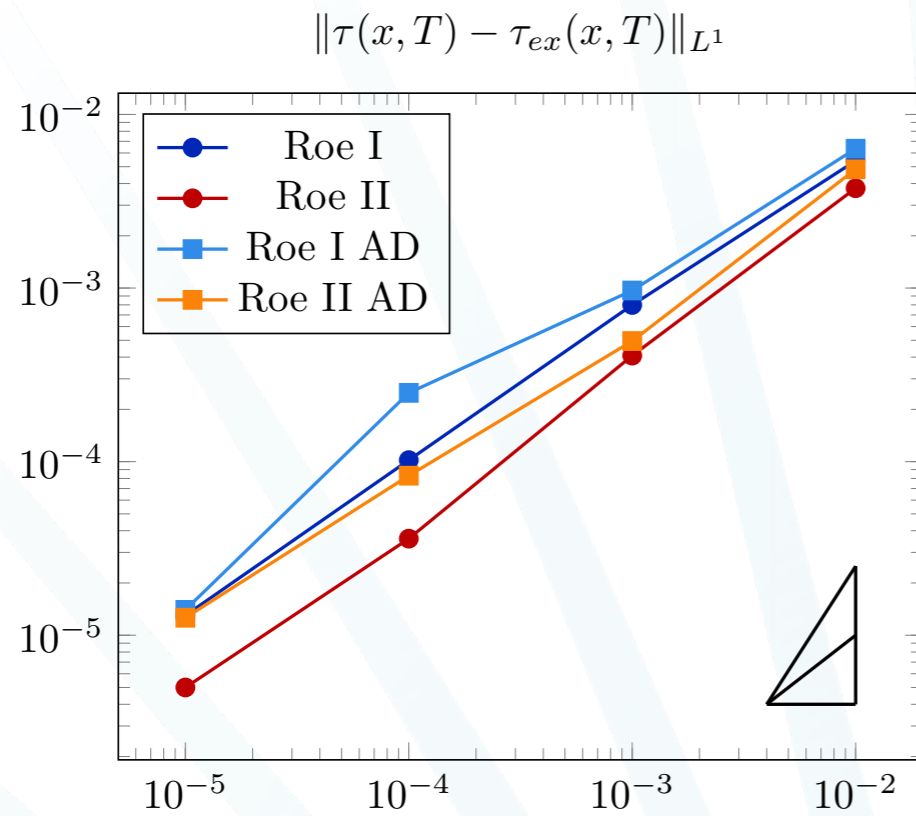




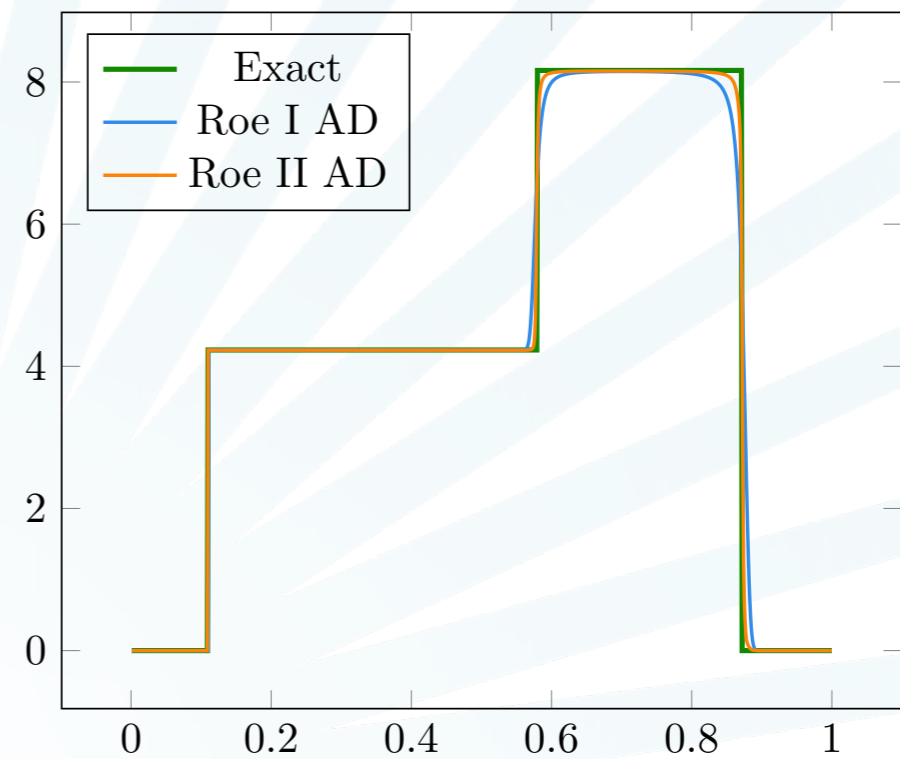
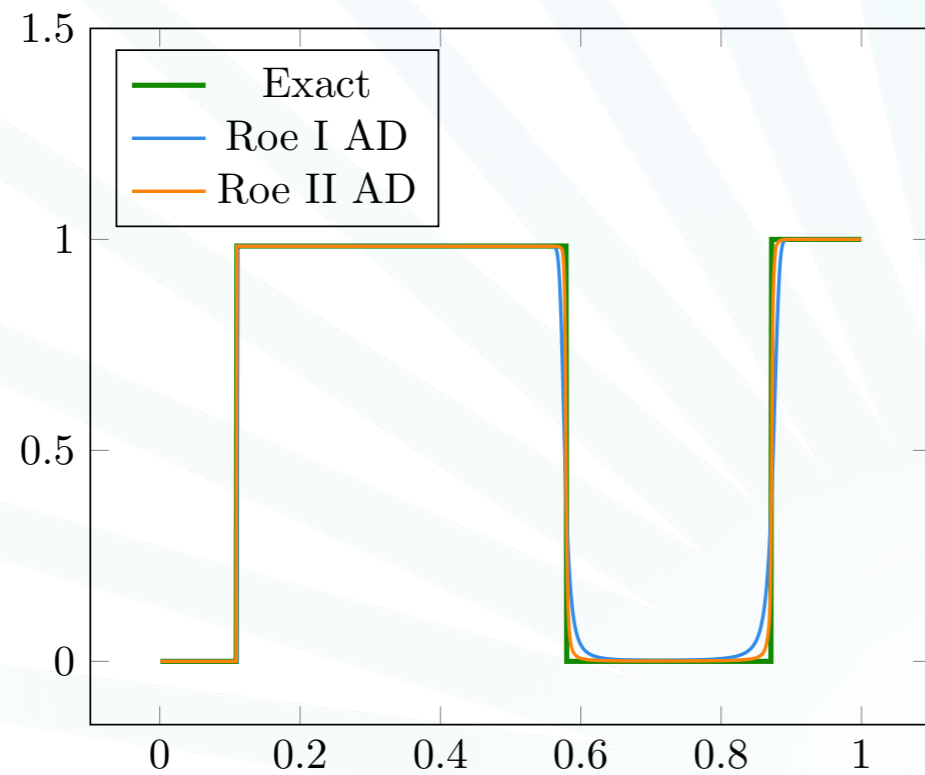
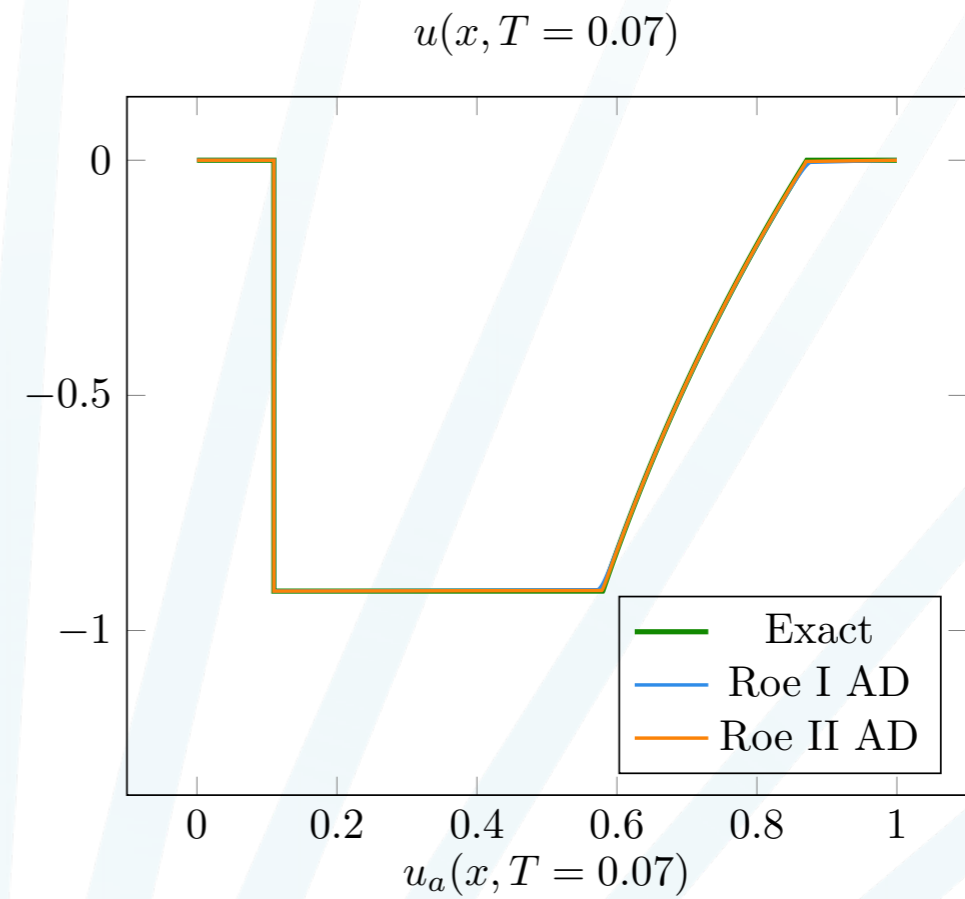
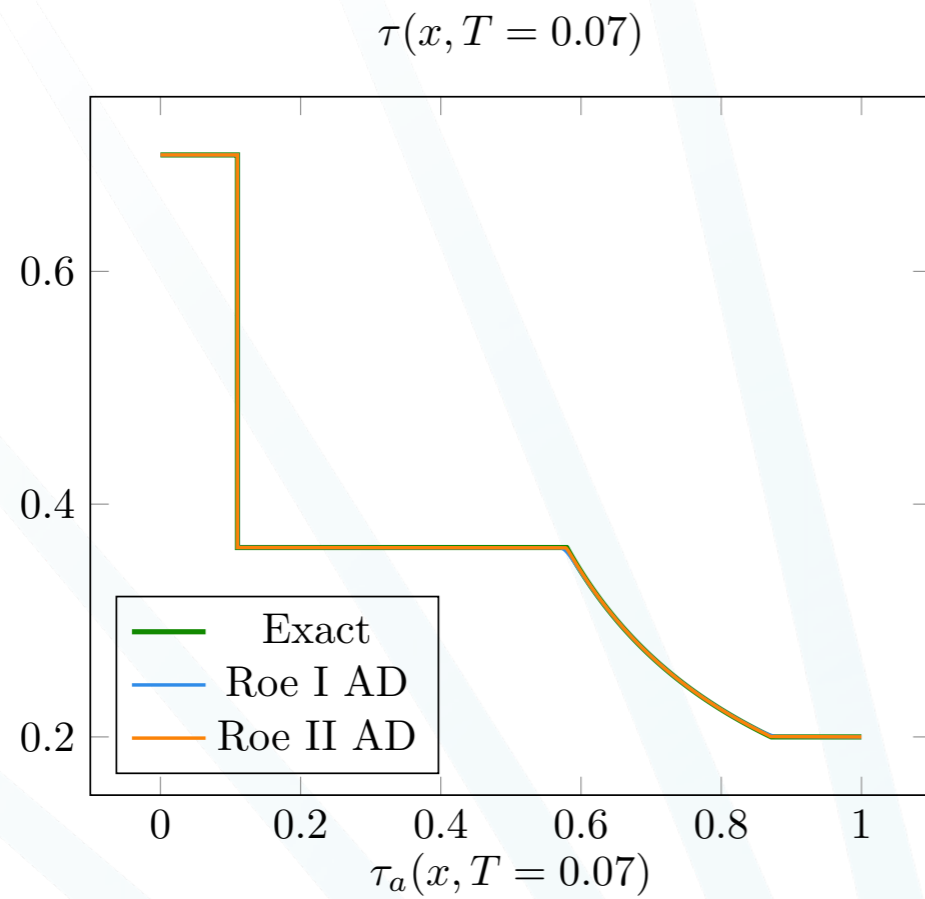
# Results



# Convergence



# Results



# Conclusion and future development

## Conclusion:

- ▶ We defined a sensitivity system valid in case of discontinuous state
- ▶ The correction term is well defined
- ▶ It is necessary to control the numerical diffusion in the shock

## Future development:

- ▶ Extension to the Euler system
- ▶ Extension to 2D
- ▶ Applications



**Thank you  
for your attention!**